

**Measurement of wakefields in  
hollow plasma channels**  
Carl A. Lindstrøm (University of Oslo)

*in collaboration with Spencer Gessner (CERN)  
presented by Erik Adli (University of Oslo)*

FACET-II Science Workshop  
Oct 18, 2017

# Measurement of wakefields in hollow plasma channels

FACET-II Science Workshop, SLAC – Oct 18, 2017

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*in collaboration with*

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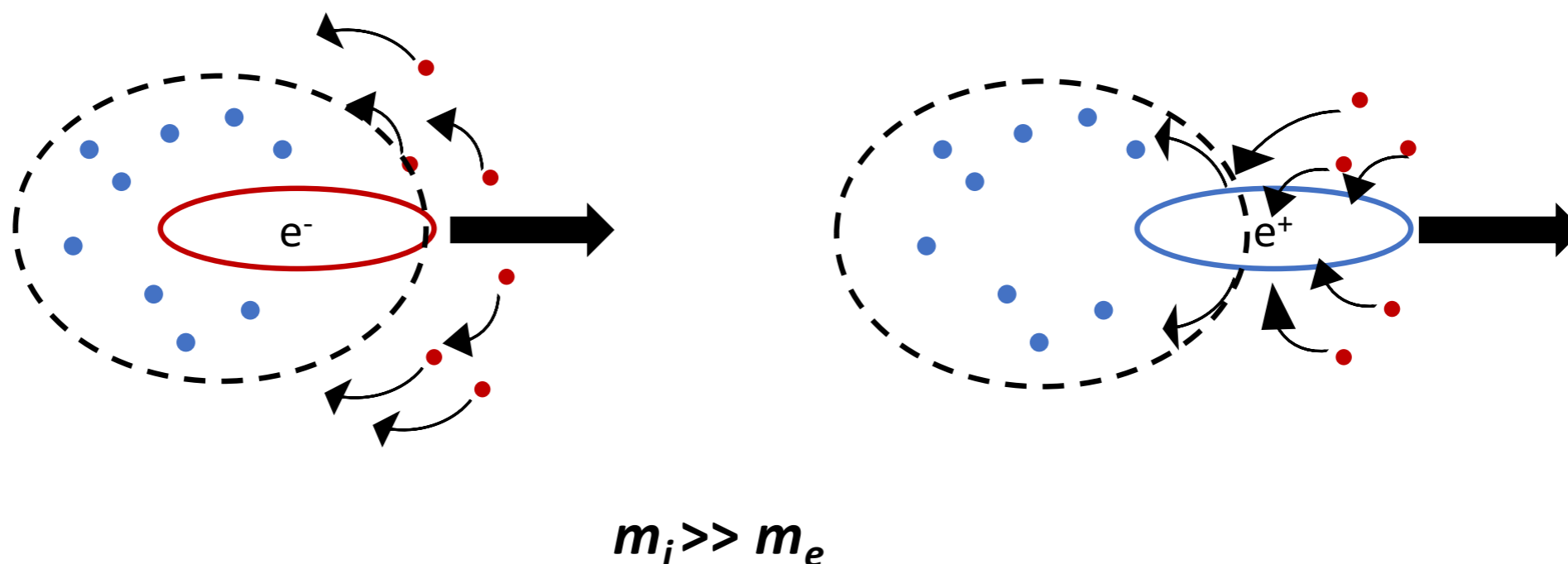


Carl A. Lindstrøm



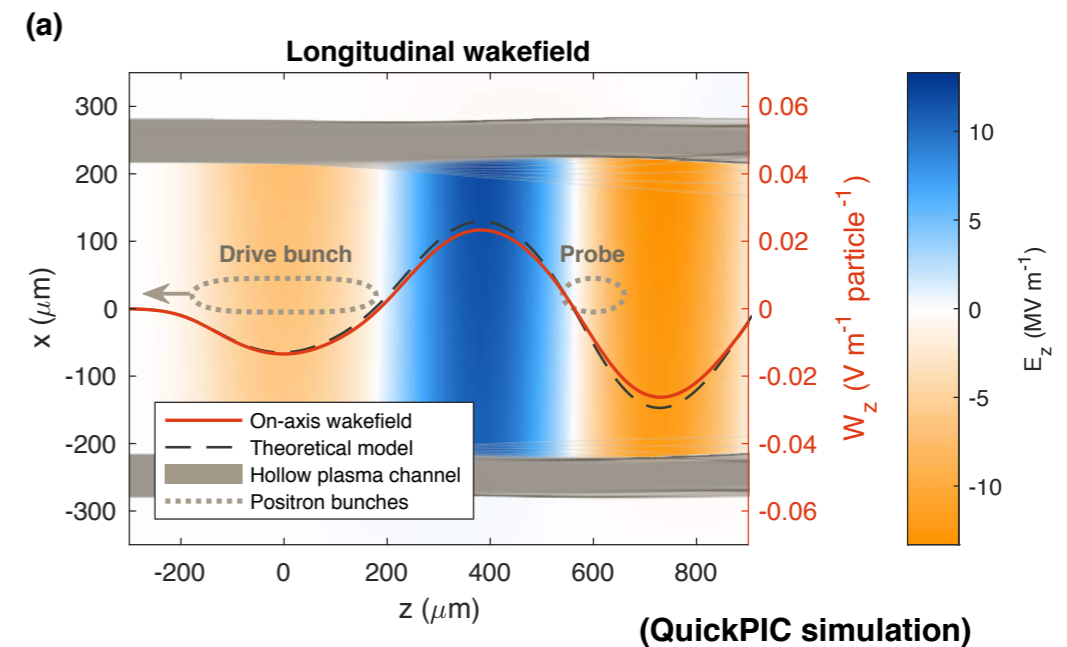
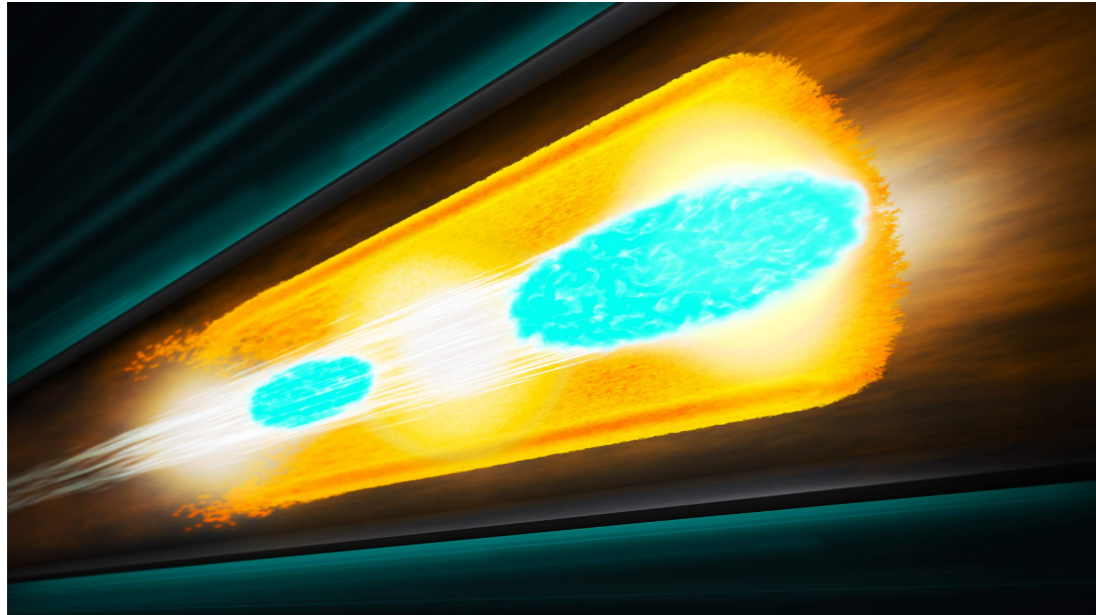
# Motivations for the hollow plasma channel

## An intrinsic charge asymmetry



- Plasmas have a problem that conventional accelerating structures do not: an intrinsic charge asymmetry.
- Even if we have a mechanism for accelerating electrons, this does not extend to positrons.

# Hollow plasma channels



- A hollow plasma channel is a proposed method to symmetrize the charge response and allow high gradient positron acceleration.
- Principle:
  - A positron bunch propagates in the centre of the hollow plasma channel
  - The channel wall is perturbed, driving an oscillating longitudinal wakefield
  - A trailing positron bunch is placed in the accelerating phase of the wakefield
- Benefit of hollow plasma channels: In principle, no focusing forces inside

# Analytical expressions for hollow channel modes exist

m=0, longitudinal mode

$$W_{z0}(z) = \frac{ek_p\chi_{\parallel}^2}{2\pi\epsilon_0 a} \frac{B_{00}(a, b)}{B_{10}(a, b)} \cos(\chi_{\parallel} k_p z) \Theta(z). \quad (1)$$

m=1, transverse dipole mode

$$W_{x1}(z) = -\frac{e\Delta x\chi_{\perp}}{\pi\epsilon_0 a^3} \frac{B_{11}(a, b)}{B_{21}(a, b)} \sin(\chi_{\perp} k_p z) \Theta(z), \quad (3)$$

S. Gessner, PhD thesis, SLAC-R-1073 (2016) ,  
earlier work by C. Schroeder (1999)

# Misalignment leads to transverse wakefields

- Drive bunches perfectly aligned to the channel axis will give zero transverse force everywhere.
- However, misaligned drive bunches will drive **strong dipole-like (transversely uniform) oscillating transverse wakefields**.
- First discussed by C. Schroeder in 1999 (“Multimode Analysis of the Hollow Plasma Channel Accelerator”).
- This leads to beam deflection and beam loss.
- This problem gets rapidly worse with stronger accelerating fields (transverse force scales faster with smaller channel radius):

**Accelerating field per particle**

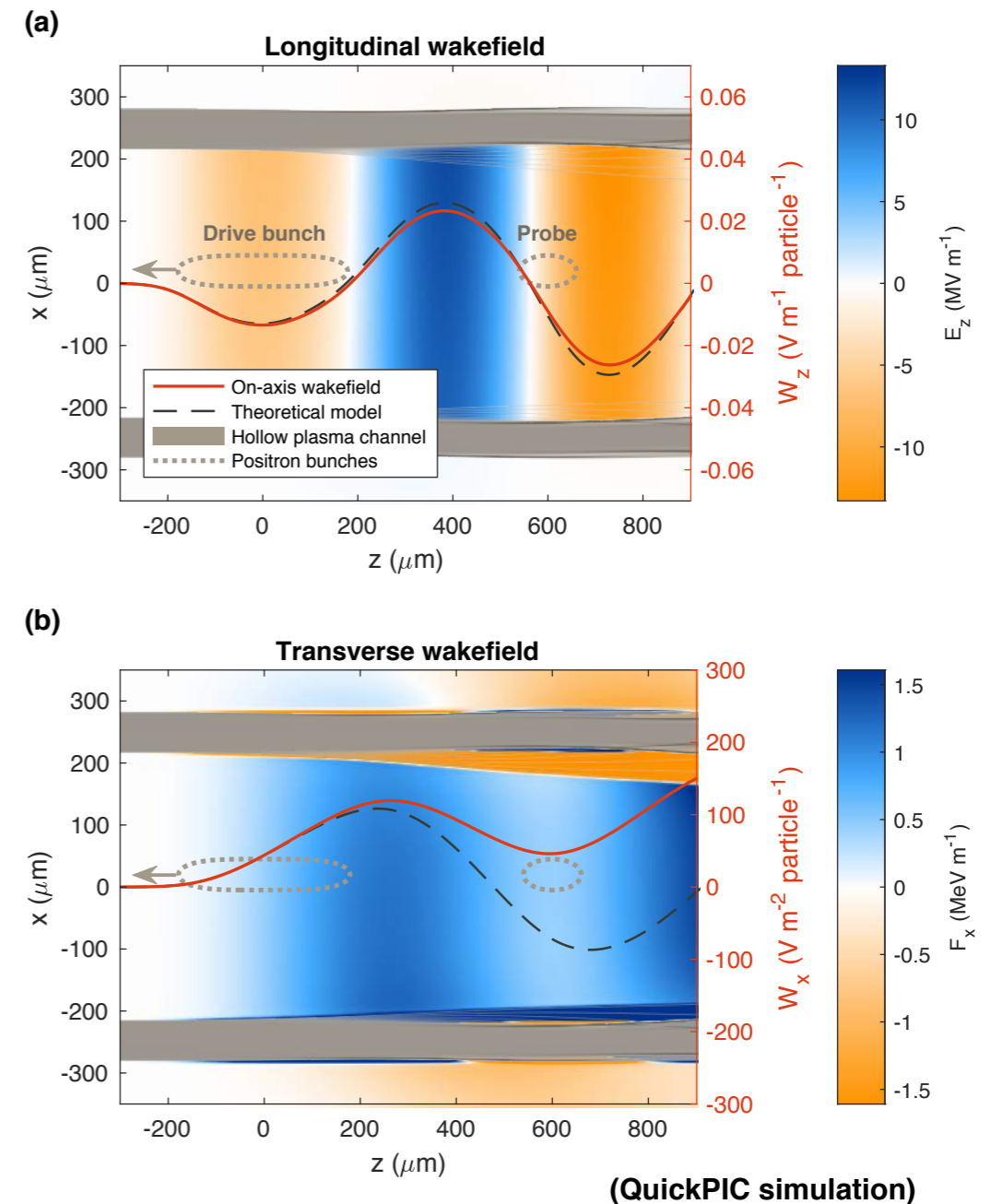
$$W_{\parallel} \propto \frac{1}{a^2}$$

**Transverse force per particle**

$$W_{\perp} \propto \frac{1}{a^{3-4}}$$

(a: channel inner radius)

- Many orders of magnitude stronger wakefields compared to CLIC
  - Hollow channel (500 μm diameter, 3x10<sup>15</sup> cm<sup>-3</sup>): ~1 000 000 V/pC/m/mm
  - CLIC (8 mm diameter, copper): ~10-100 V/pC/m/mm

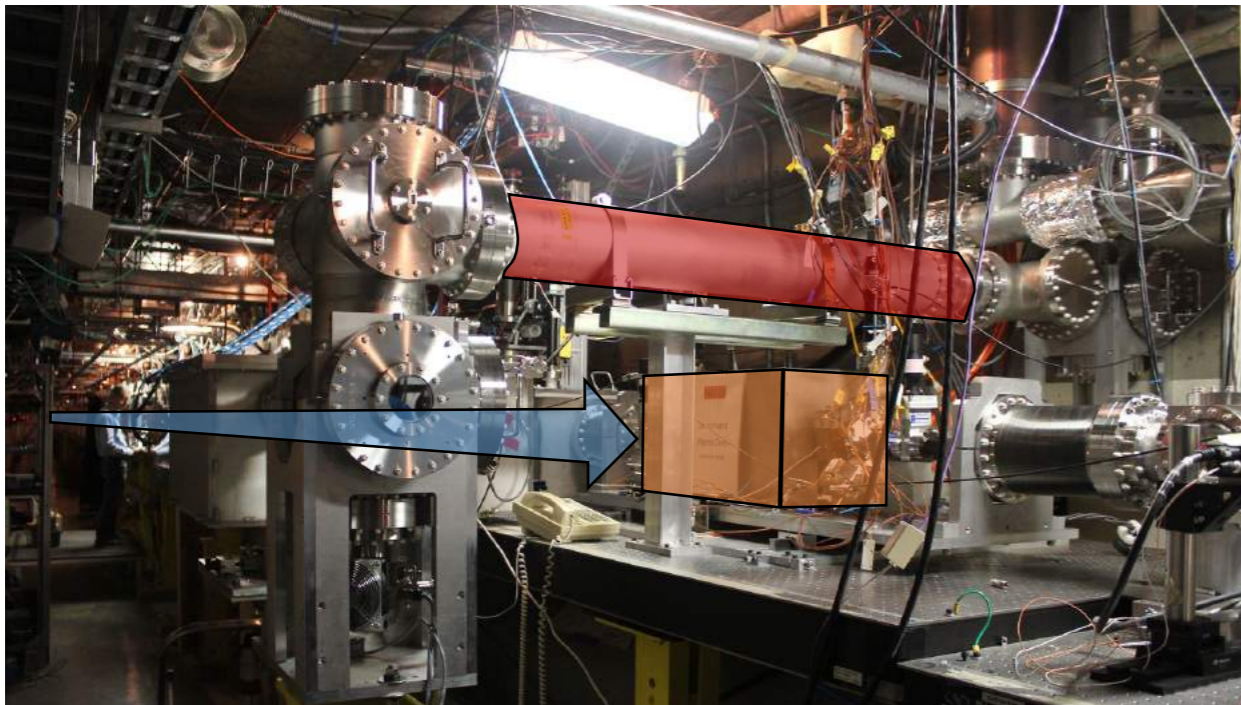




# Experiments at FACET

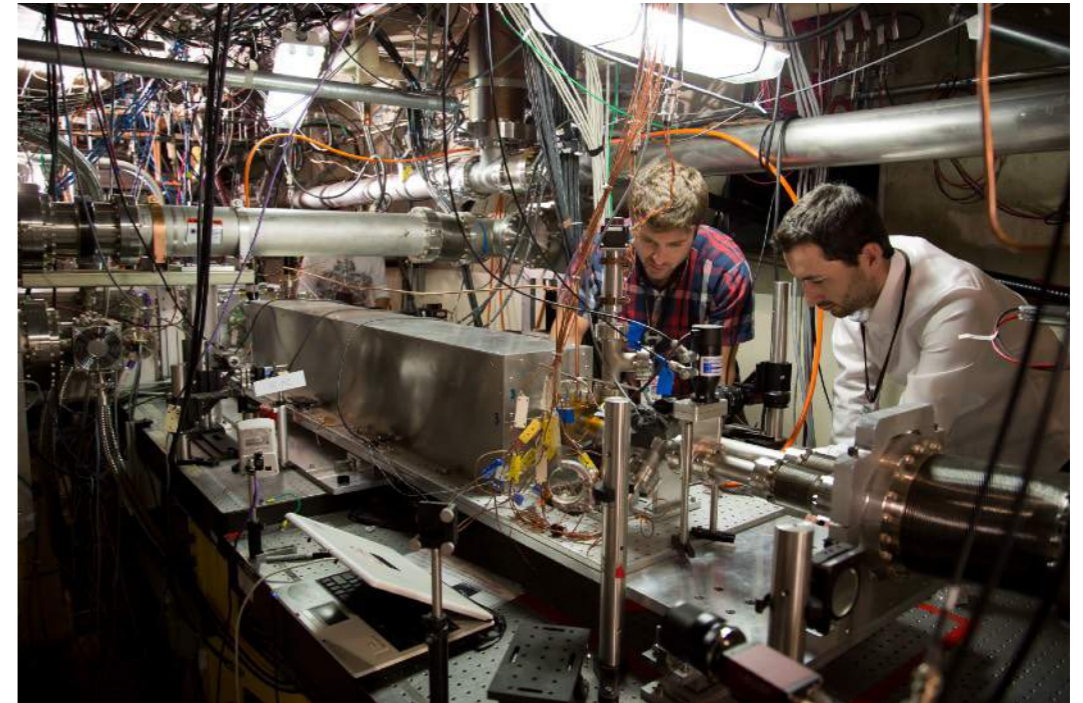
## The E225 experiment

- FACET hosted the dedicated hollow plasma channel E225 experiment, lead by Spencer Gessner
- The main aim was to demonstrate positron acceleration in a hollow channel, but also to investigate transverse wakefields



**FACET experimental area, showing positron beam (blue), ionising laser (red) and lithium vapor oven (orange).**

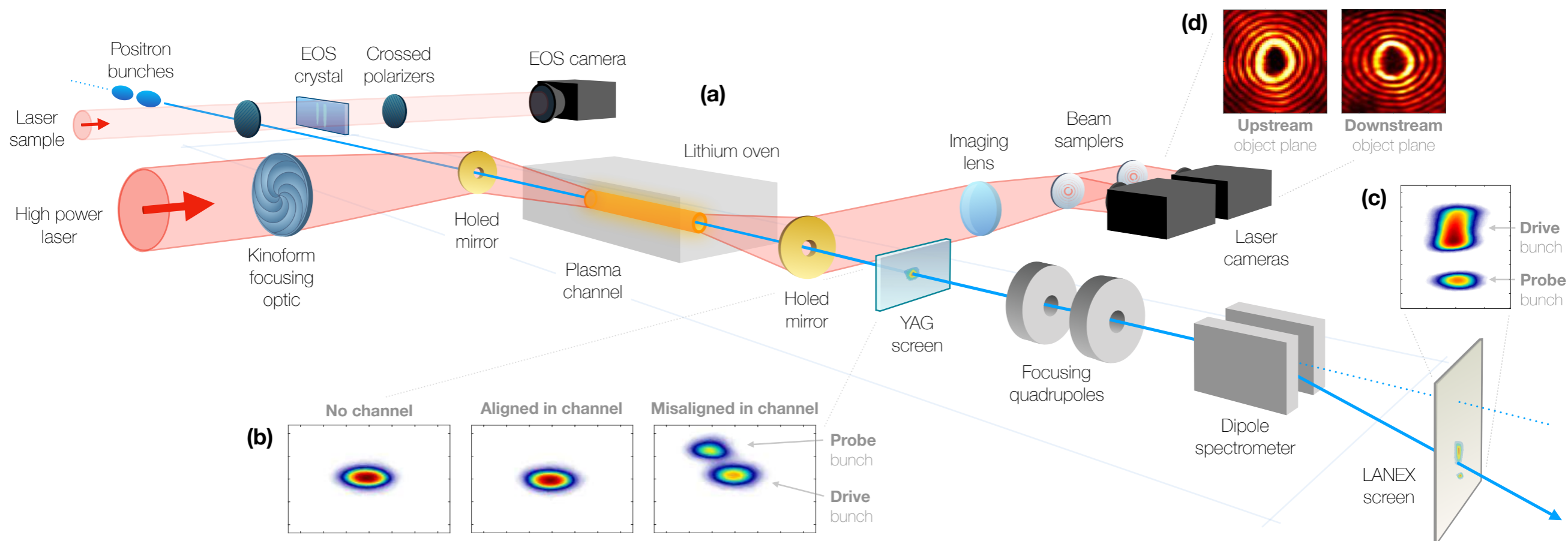
*Image source: Spencer Gessner*



**Spencer Gessner (left) and Sebastien Corde (right) at FACET.**

*Image source: SLAC National Accelerator Laboratory*

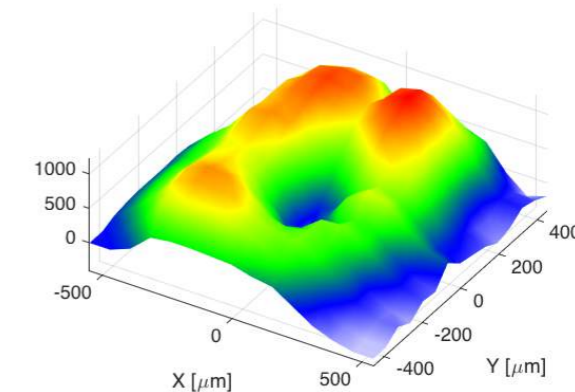
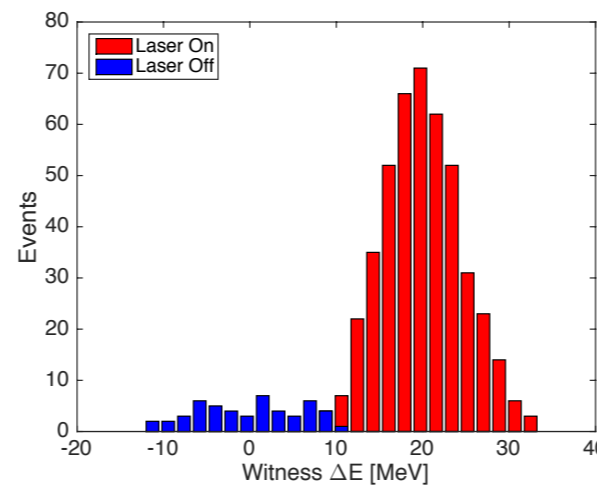
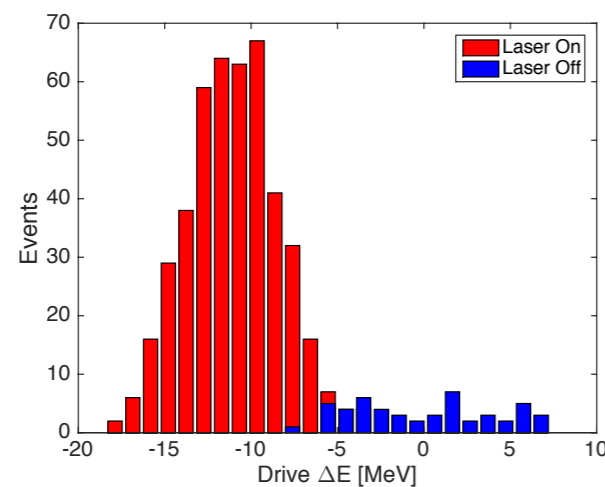
# E225 – Experimental setup



- The SLAC linac provided **two 20 GeV bunches**, made from one bunch using a beam notching device.
- The FACET laser (up to 10 TW, 60 fs pulses) was adjusted down to ensure **no ionisation in the channel**.
- A lithium oven was set to give a neutral gas density of  $3 \times 10^{16} \text{ cm}^{-3}$  (but was necessarily fully ionized).

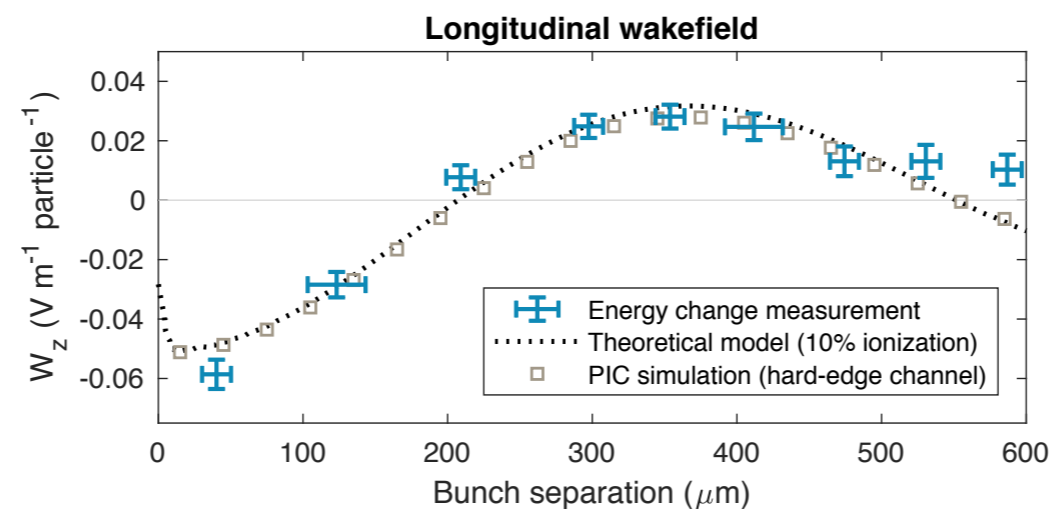
# Positrons successfully accelerated in a hollow channel

- Clear evidence of energy gain for the positron witness bunch, while there is energy loss for the positron drive bunch.



Reconstruction of the plasma channel based on kick measurements

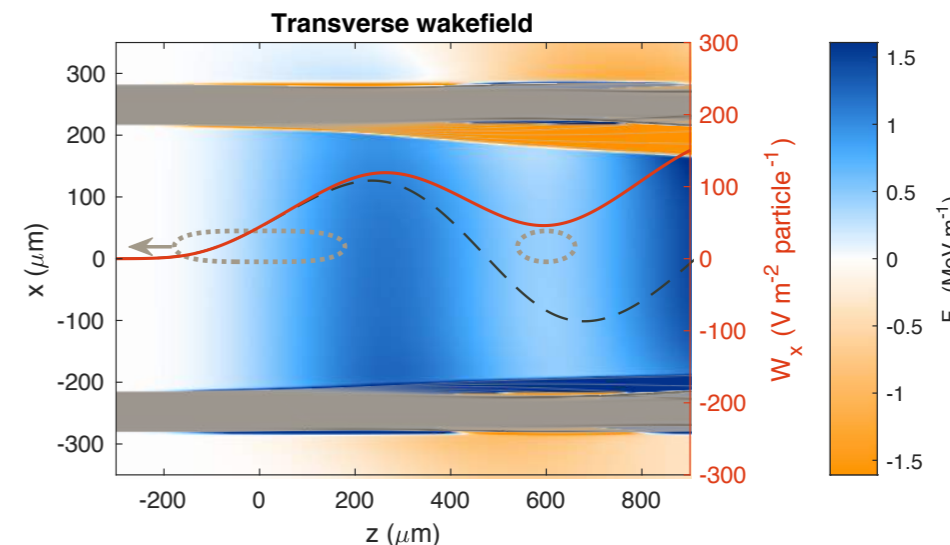
- A scan of bunch separations shows the energy gain (or loss) depends on the phase.



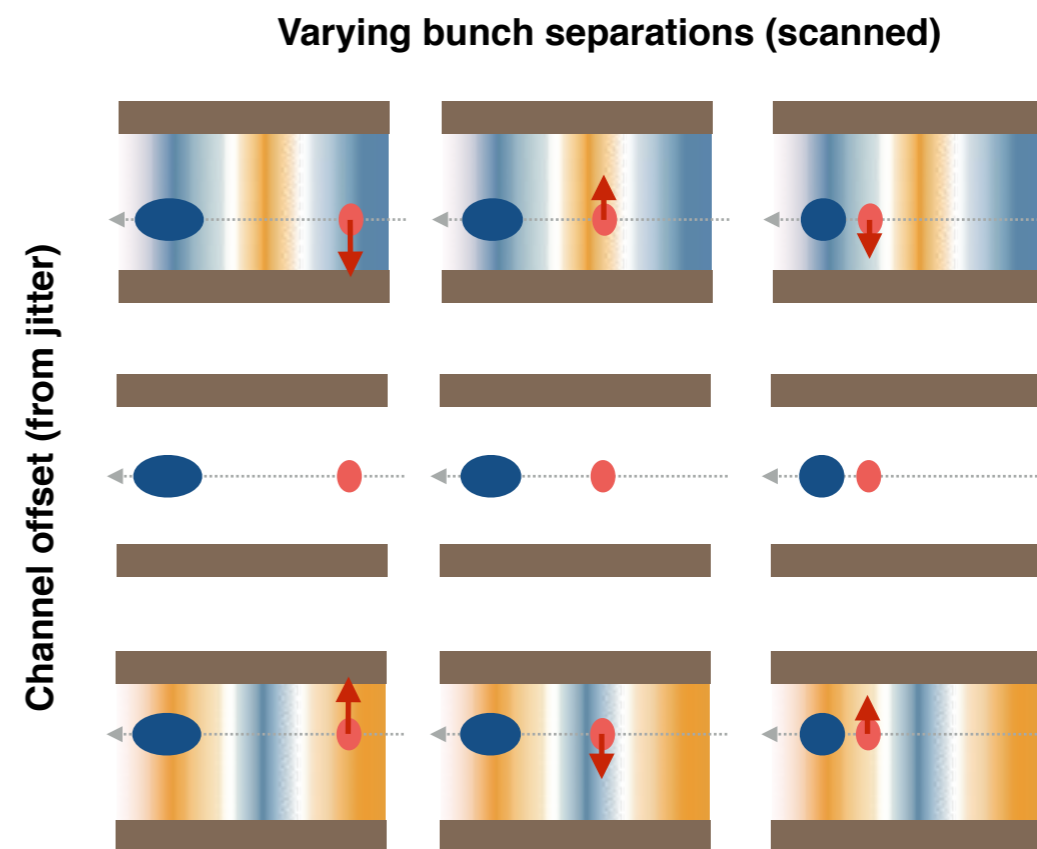
# The transverse wakefield experiment

- Our **goal** was to **measure the how the transverse wakefield varied longitudinally**.
- The probe bunch observing the wakefield is deflected angularly (kicked) when the channel and the drive bunch are relatively offset.
- The experiment performed was:
  - Transverse channel offsets for various bunch separations**
    - The channel (250  $\mu\text{m}$  radius) was offset by transverse laser jitter (20-40  $\mu\text{m}$  rms)
    - The bunch separation was varied by stretching the bunch and adjusting the notching device.
- Diagnostics:
  - Laser offset imaged downstream (**laser cameras**).
  - Probe kick measured on a **spectrometer** (in the non-dispersed plane).
  - Bunch separation measured using an **electro-optical sampler**.

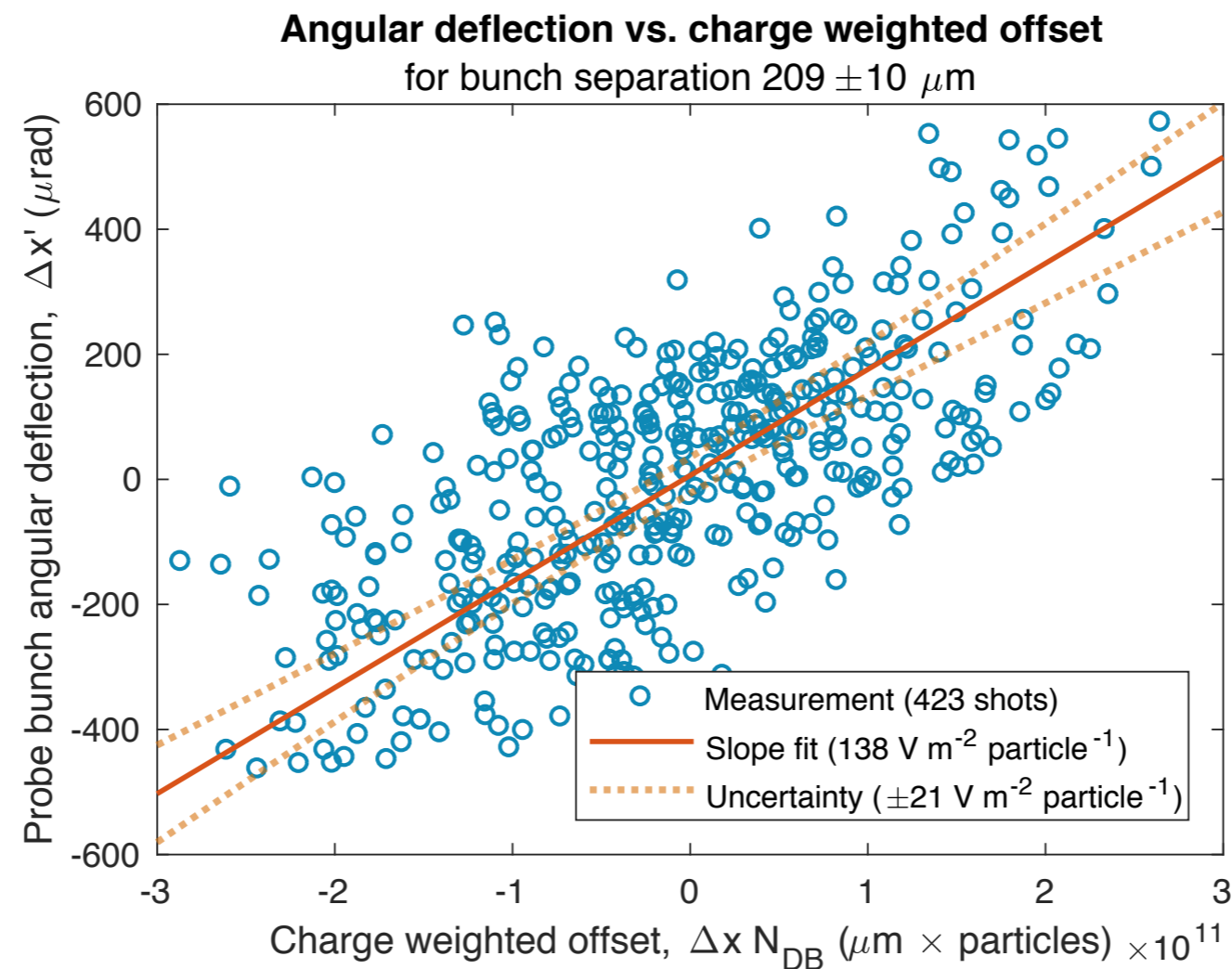
Prediction:



Experiment (2D “scan”):



## Observed data (deflection vs. channel offset)



- For each bunch separation, a correlation between channel offset and probe bunch angular deflection was observed.
- The **slope of this correlation is proportional to the transverse wakefield per offset** at the z-location of the probe bunch.

## Another independent measurement

- An independent measurement is beneficial (due to high complexity).
- It is possible to estimate the transverse wakefield per offset from the measured longitudinal wakefield, via the Panofsky-Wenzel theorem and the linear model.

*Panofsky-Wenzel theorem:*

$$\frac{\partial W_x}{\partial z} = \frac{\partial W_z}{\partial x}$$

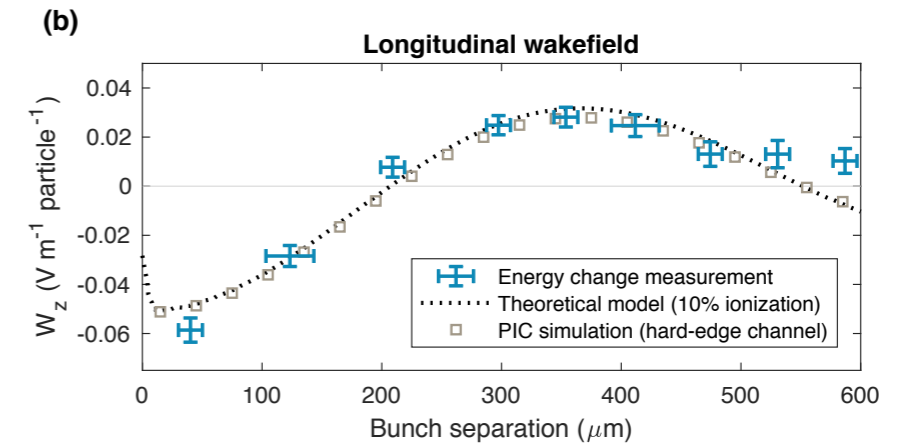
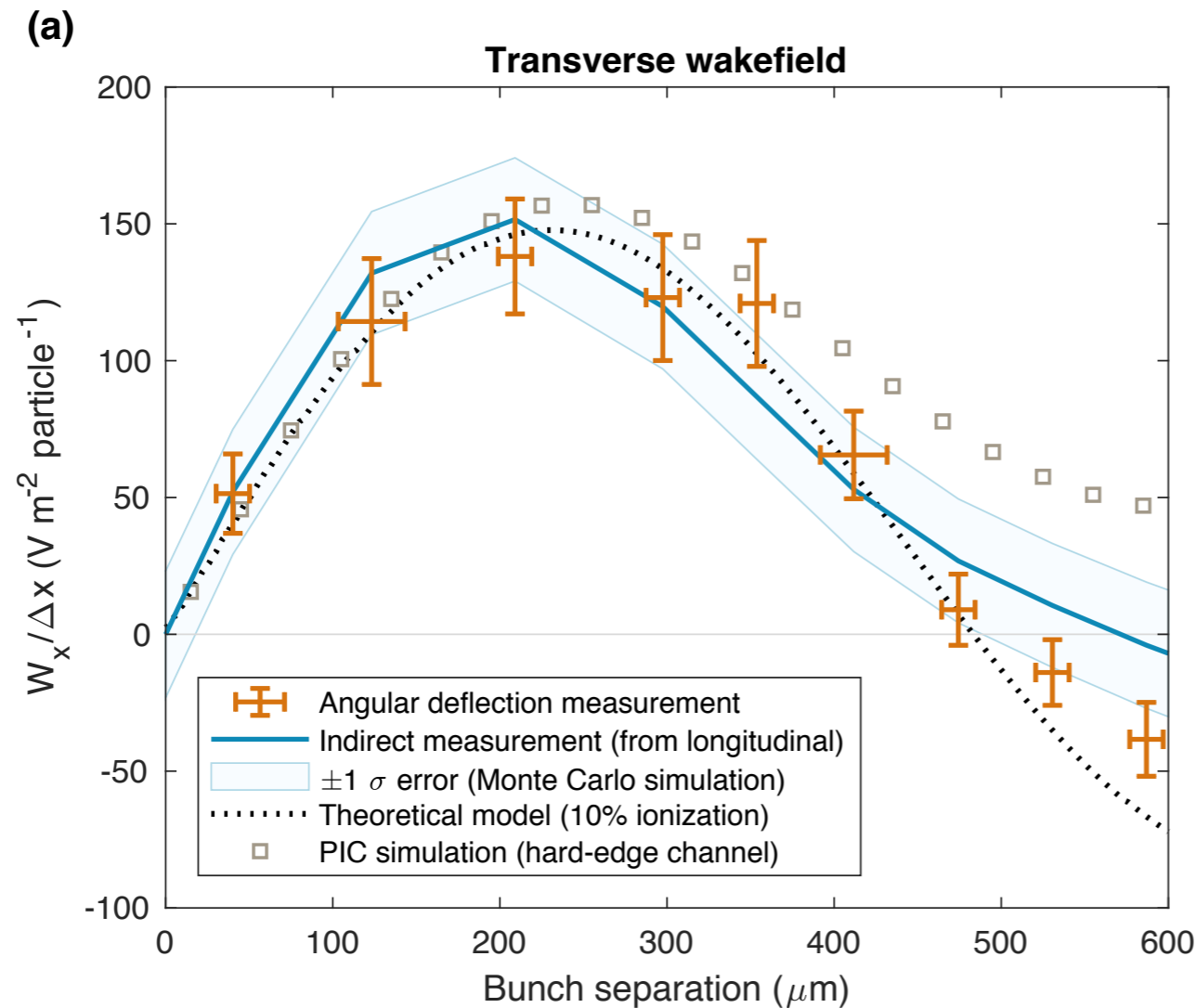
Integrate (++)  
→

*Estimate of transverse from longitudinal wakefield:*

$$\frac{W_x(z)}{\Delta x} \approx -\frac{\kappa(a, b)}{a^2} \int_0^z W_z(z') dz' \quad \text{where} \quad \kappa(a, b) = \frac{4\chi_{\perp}^2 - 2}{\chi_{\parallel}^2 - 1}$$

- Not perfect: Assumes linear model, breaks down far behind the drive bunch.
- Provides verification of numerical calibrations, etc.
- The longitudinal wakefield was measured by the energy change of the probe bunch (on a spectrometer).

# Final experimental results

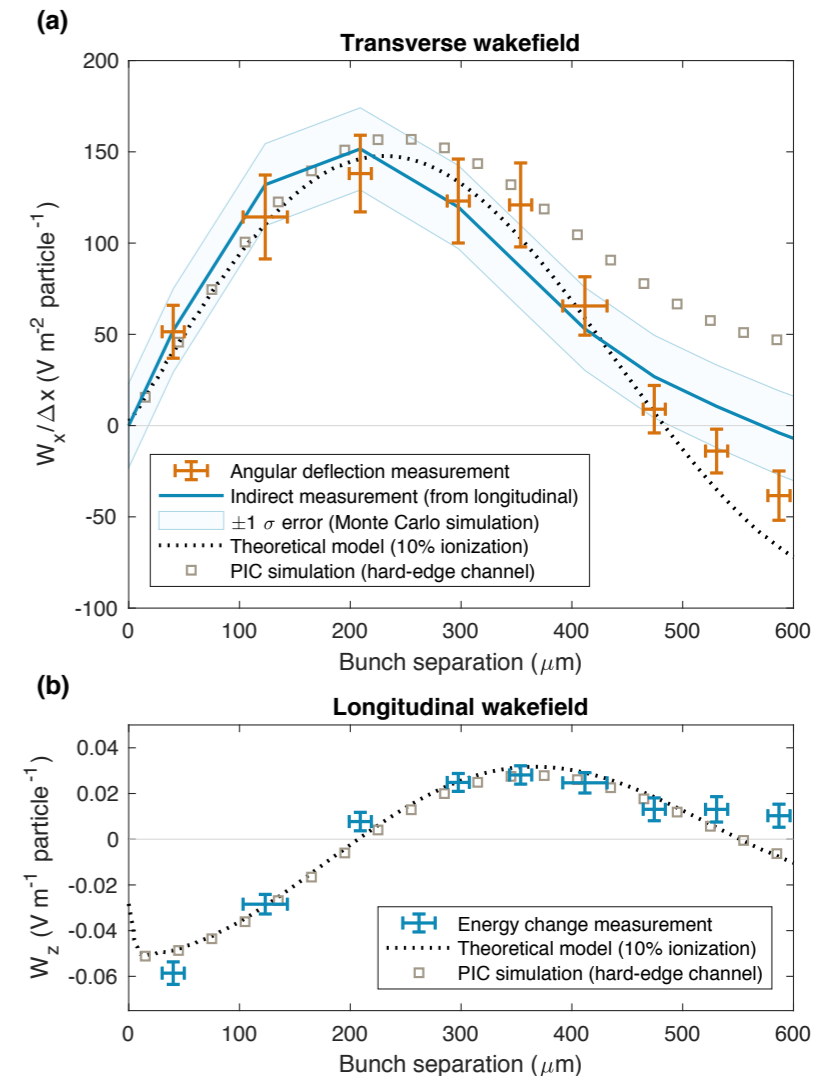


- Plasma density determined by a wavelength fit (10% ionization =  $3 \times 10^{15} cm^{-3}$ )
- Good fit, largely consistent with theory. Some discrepancy at larger separations.



## Implications

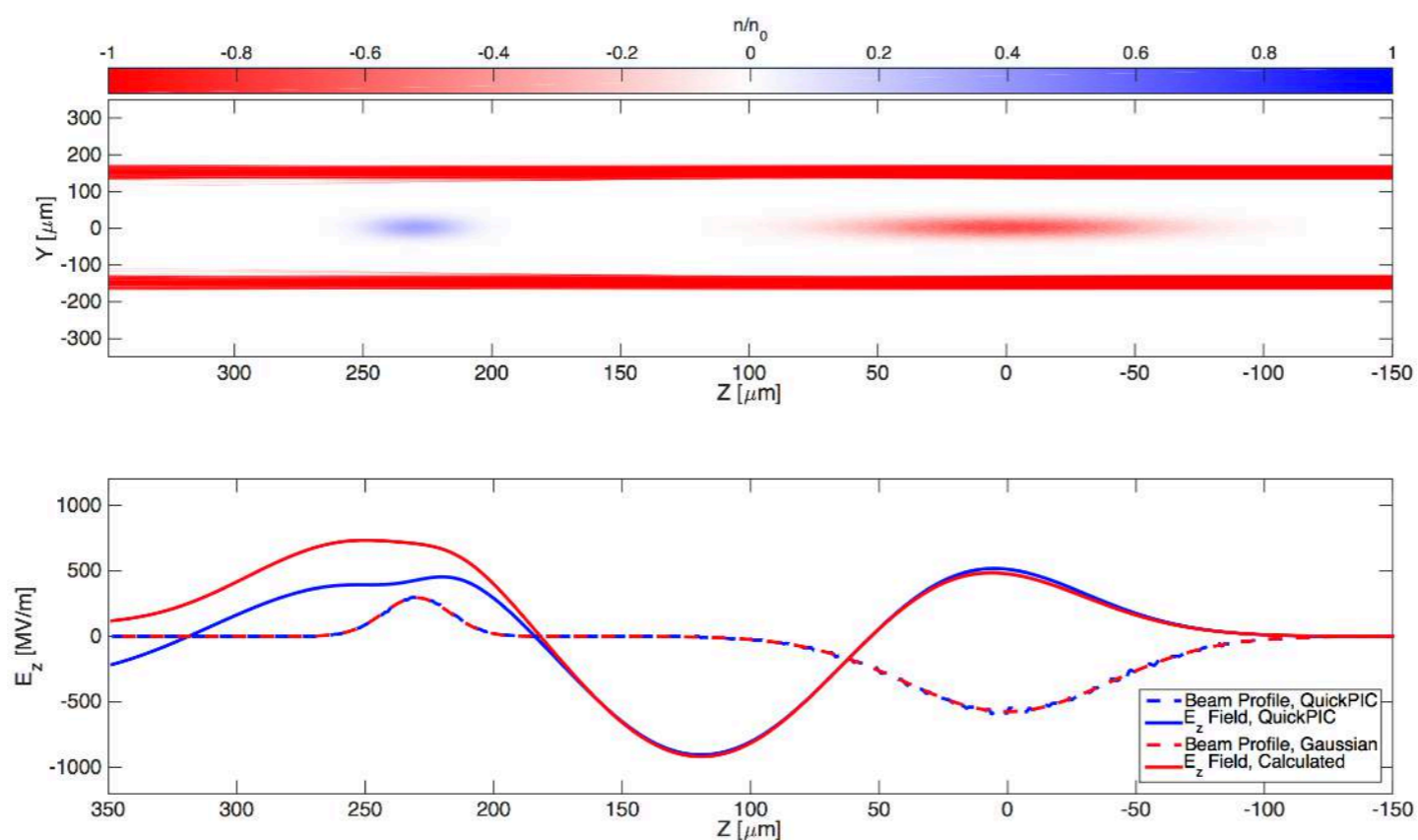
- **Overall, the measurement agrees with the theoretical models.**
- Simulation-based parameter scans indicate that the discrepancy at large separations can possibly be explained by using a more complex radial plasma shape (not possible to exclude with our diagnostics).
- Implication:
  - There is indeed a strong transverse wakefield, as expected.
  - This needs to be mitigated for the hollow channel to be useful.
- Submitting these results to Physical Review Letter.



# Future directions

## Positron acceleration should be electron driven

- Wall plug-to-beam efficiency is key to high energy colliders
- Therefore, ideally a plasma-based linear collider is electron driven (or proton driven) because positrons are energy intensive to produce.



Electron driven positron acceleration in a hollow plasma channel.

Image source: **S. Gessner thesis, SLAC-R-1073**

Two big challenges for hollow plasma channels

Problem #1 (fundamental)

**Suppressing the transverse wakefield**

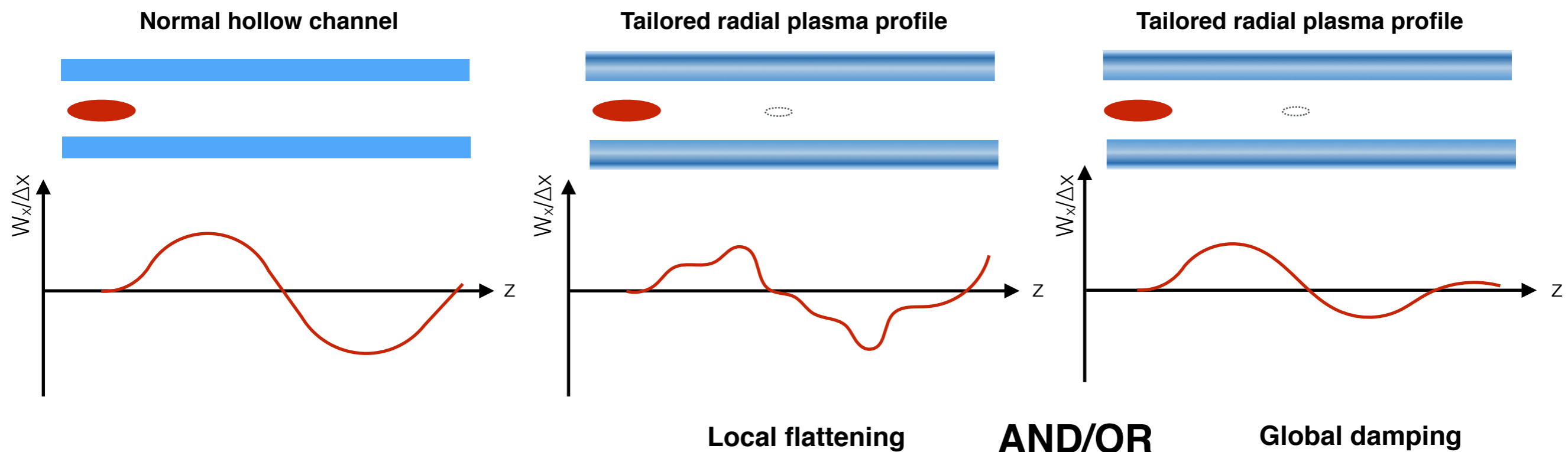
Problem #2 (technical)

**Creating an on-axis vacuum**

# Suppressing the transverse wakefield (one interesting pathway)

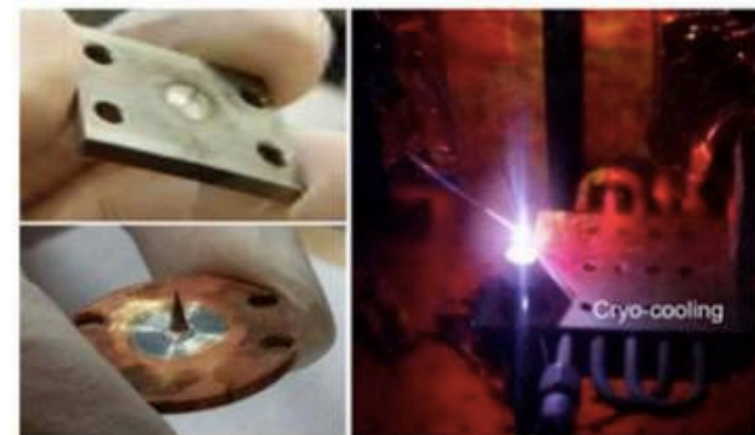
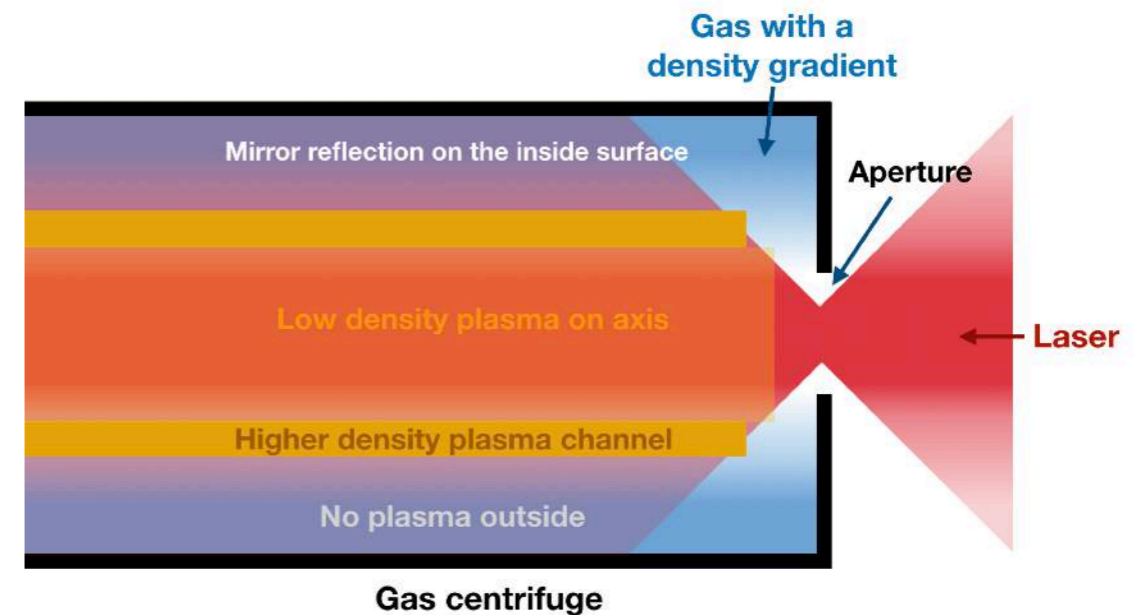
- The wakefields are determined in part by the radial plasma profile  $n(r)$ .
- Speculation: A suitably tailored radial profile may damp (locally or globally) the transverse wakefield, while sustaining a non-zero accelerating field

(ref. G. Shvets "Excitation of Accelerating Wakefields in Inhomogeneous Plasmas", 1996)



## Creating a vacuum on axis

- Eventually, we need to have a vacuum on axis, to avoid beam ionisation.
- **Centrifuge technique**, where the gas density is approximately exponentially decaying towards the axis.
- **Cryo-cooled gas cluster** technique (used for corrugated plasma channels by H. Milchberg)
- These ideas can potentially be tested in the laser labs at UCLA or UC Boulder.



**Image source:**  
H. Milchberg (Uni Maryland), EAAC2017 talk

## Conclusion

- Hollow channels are promising, supporting very strong longitudinal wakefields.
- However, they also support very strong well as transverse wakefields (leading to beam loss)
- Positrons were accelerated in a hollow plasma channel!
- The transverse wakefield was measured experimentally, and found largely consistent with theory.
- Suppression mechanisms for the transverse wakefield is key to the survival of the hollow channel.

## Ideas for FACET-II hollow channel experiments

- Radial tailoring of the hollow channel profile (laser shapes, time delays, etc.)
- On-axis vacuum (centrifuges, cryo-cluster flow, etc.)
- Electron-driven hollow channel positron acceleration.

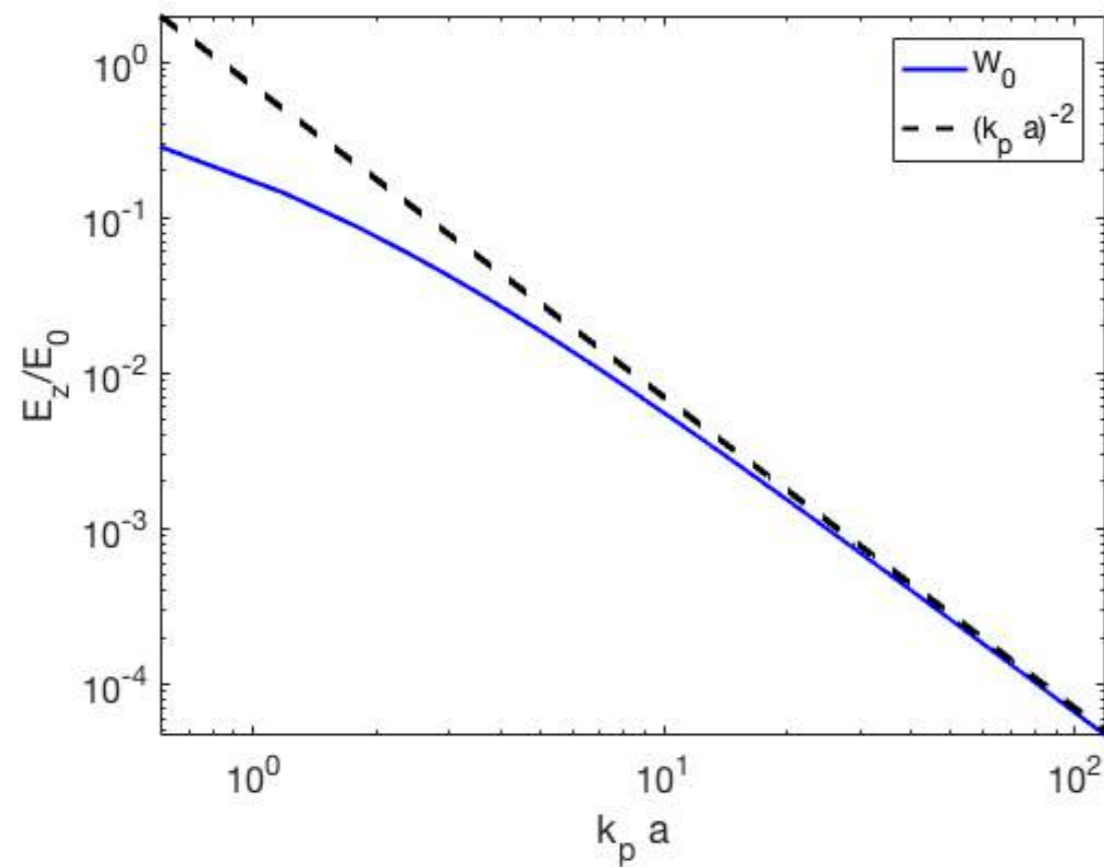
**Thank you for your attention!**



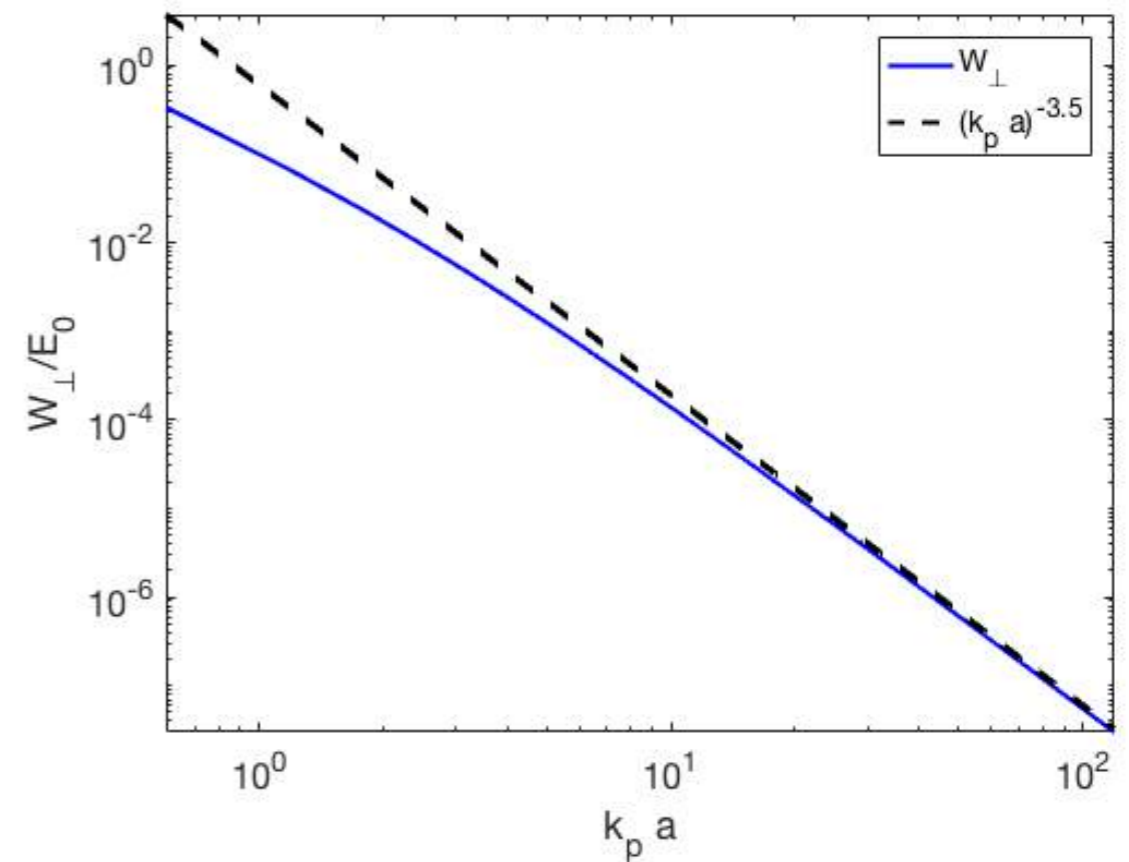


# Dependence on channel radius

## Longitudinal wakefield



## Transverse wakefield



by Spencer Gessner