

Teravolt-per-meter plasma wakefields from low-charge, femtosecond electron beams

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Brightness Begets Brightness: Electrons and Photons

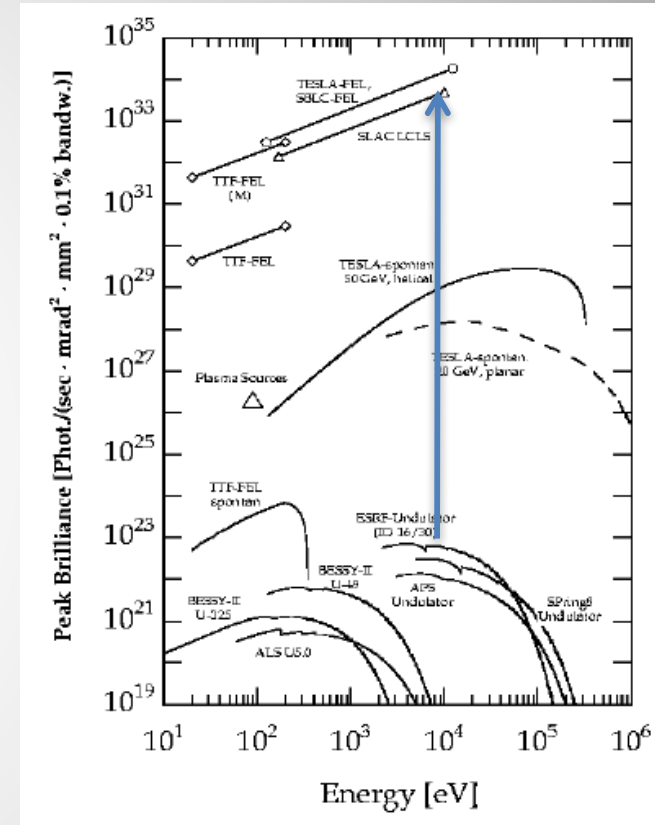
- Light source revolution due to e- beam improvements

$$B_e = \frac{2I}{\epsilon_n^2} \rightarrow 2 \times 10^{14} \text{ A/m}^2$$

- Two orders of magnitude produce qualitatively new *light source*, the X-ray Free-electron Laser
 - Intense cold beam; instability

$$E_{rad} \propto \exp(z / L_g); L_g \propto B_e^{-1/3}$$

- 8 orders of magnitude photon brightness
- For needed e-beam, must compress to **fs scale**
- Ultra-bright, Å, coherent, fsec light source
 - X-ray FEL: many orders of magnitude leap forward

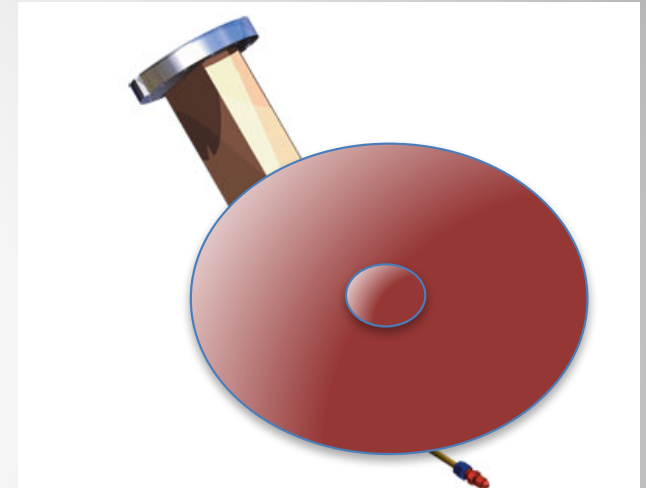


Ultra-short XFEL pulses: motivation

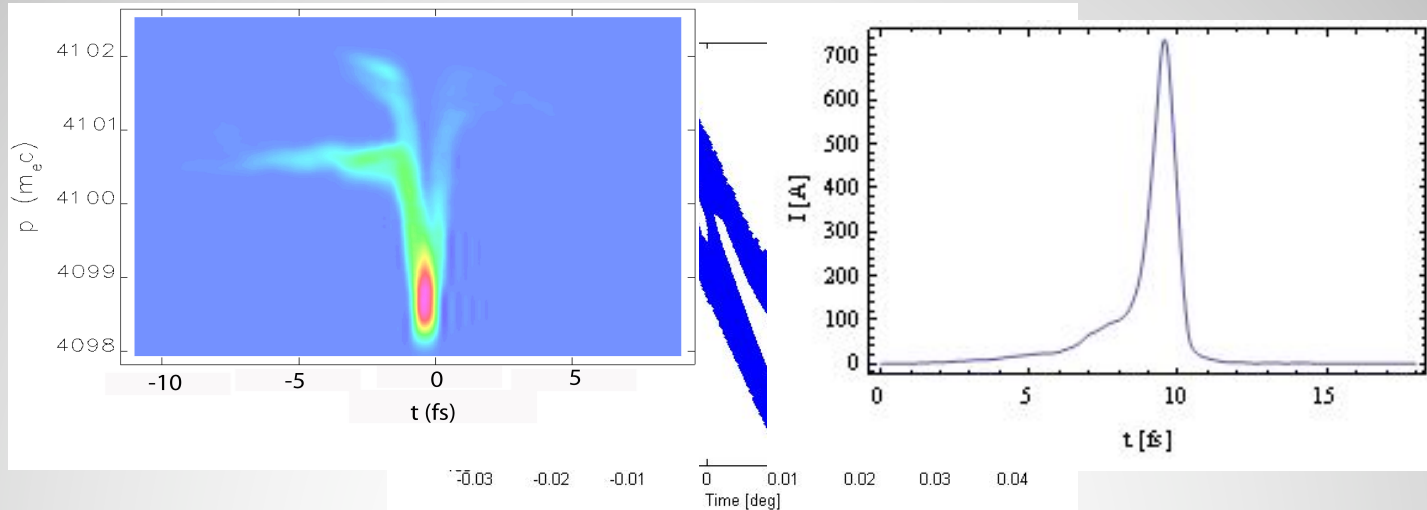
- Tool at atomic *electron* spatio-temporal scales
 - Angstroms-nanometers (\sim Bohr radius)
 - **Femtoseconds** (e^- motion, Bohr period; femto-chemistry, etc.)
- 100 fs accessible with standard approaches
- Promising path: ultra-short, low Q electron beam
 - Myriad of advantages in FEL *and* beam physics
 - Mitigate collective effects dramatically
 - Robust in application: XFEL, coherent optical/IR source
- Can also use microbunching...
- ***Spin-off* to ultra-high field PWFA!**

Beam physics: from plasma to plasma

- Beam at lower energy is single component relativistic *plasma*
- Preserve optimized dynamics: change Q , keeping plasma frequency (n , aspect ratio) same
- Dimensions scale as $\sigma_i \propto Q^{1/3}$
 - Shorter beam, easier to *compress*
 - Big emittance reduction, easy to *focus*
 - Result: ultra-high *brightness* beam



Ultra-short pulses at SPARX (LNF)



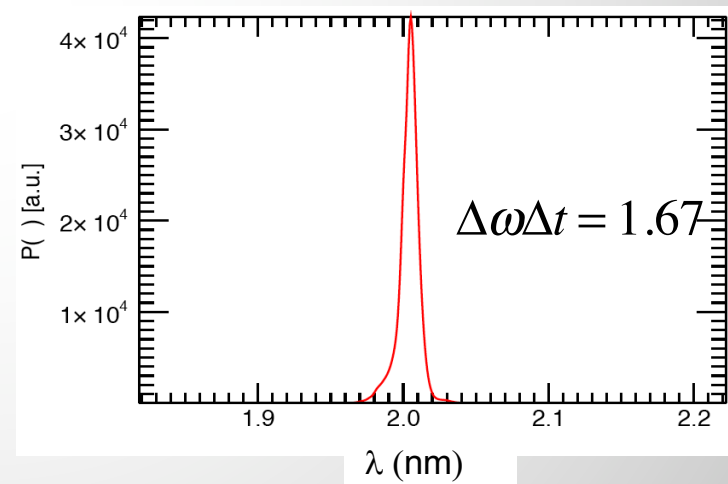
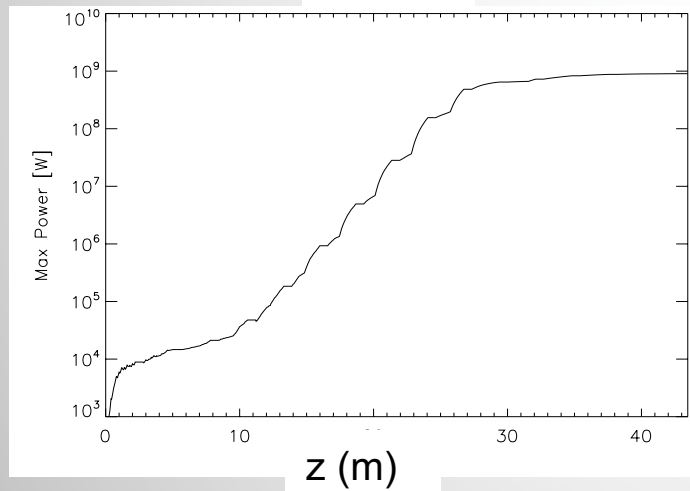
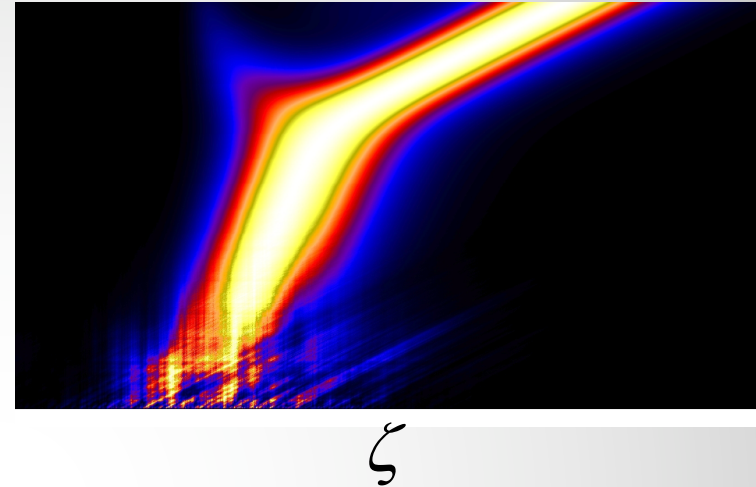
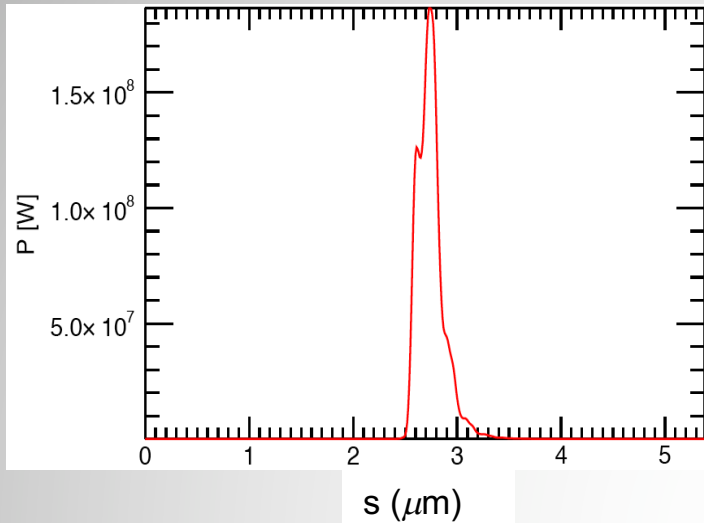
Final longitudinal phase space $\sigma_z = 4.7 \mu\text{m}$ after velocity bunching

Final current profile

- Chicanes bunching after velocity bunching
- Use ~ 1 pC beam for single spike
 - SS: cooperation length=bunch length
- Short, low emittance beam at final energy 2.1 GeV
 - $\varepsilon_{nx} \cong 7.5 \times 10^{-8}$ m-rad $\sigma_t \cong 600$ attoseconds(!)
- Very high final brightness
 - 2 orders of magnitude!

$$B = 2 \times 10^{17} \text{ A/m}^2$$

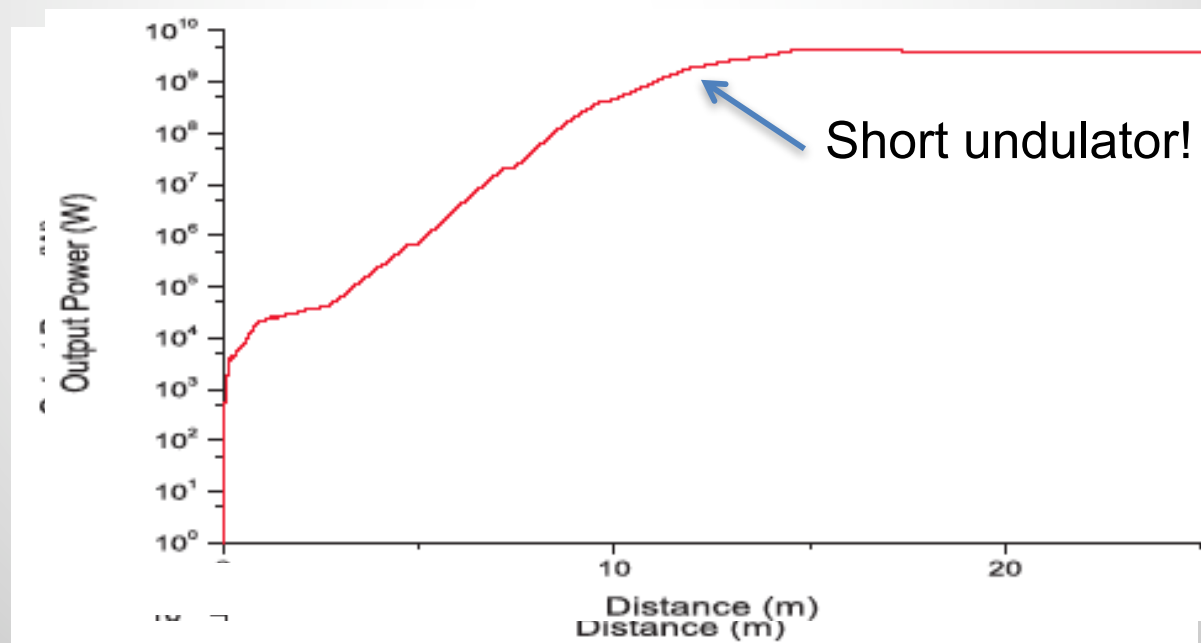
Single Spike X-ray FEL



- Single spike, > 1 GW peak power
- 480 *attosecond* rms pulse at 2 nm
- 1st time in X-ray regime

Example: LCLS w/sub-fs pulse

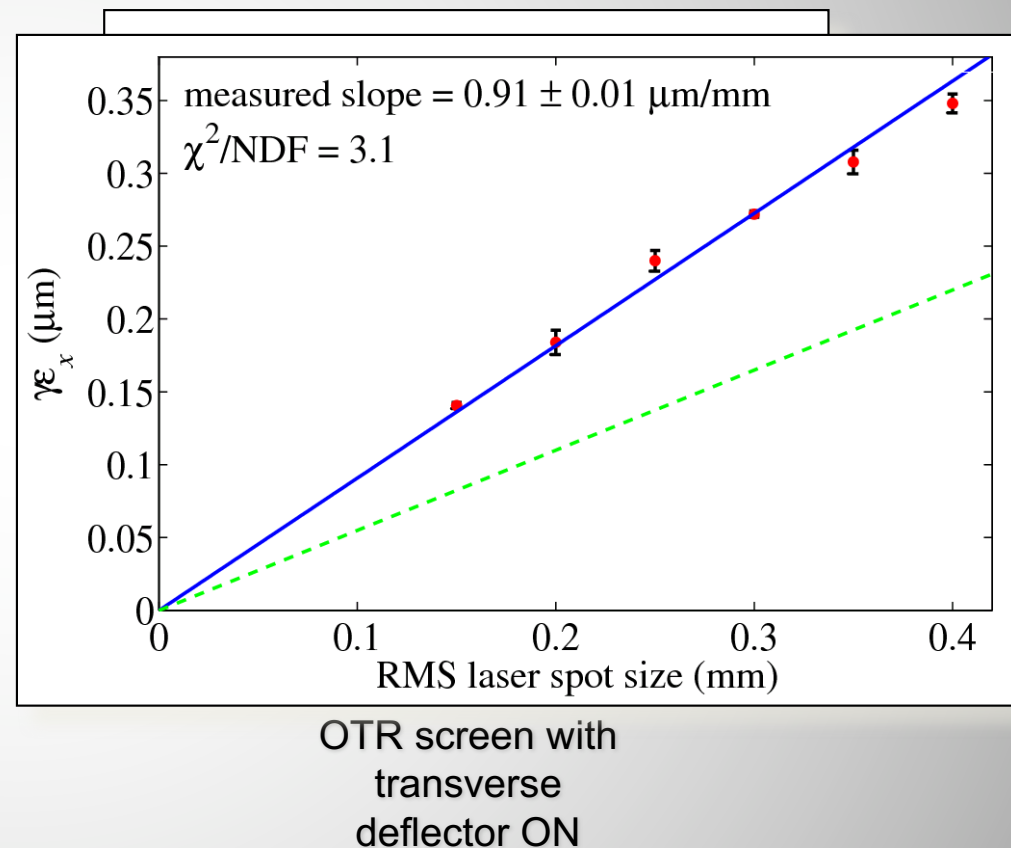
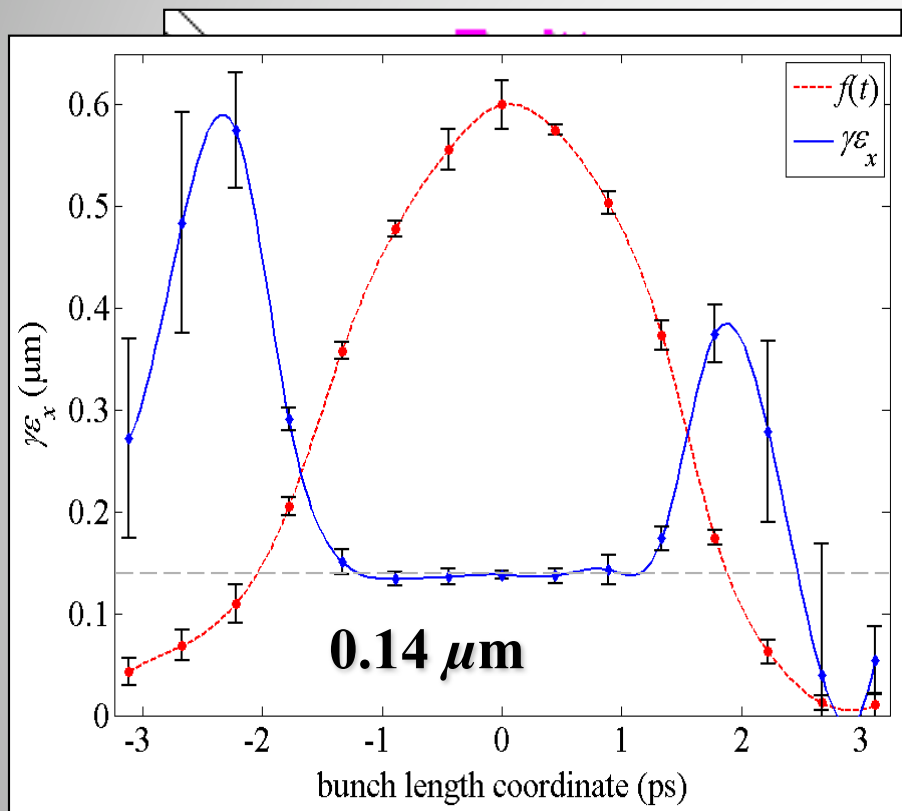
- Use even shorter 0.25 pC beam, *150 as* pulse
 - Single spike w/standard LCLS undulator
- Obtain ultra-compact “LCLS” at 4.3 GeV
- Extend energy reach to 83 keV (0.15 Å)



Gain evolution for 1.5 Å, 4.3 GeV, (0.25 pC)
Gain evolution at 13.6 GeV, 0.15 Å, 83 keV FEL

High interest in FEL community... low-Q explorations at LCLS

Low emittances at LCLS with 20 pC. Diagnostic limited



Emittance near calculated thermal emittance limit

20 pC, 135 MeV, 0.6-mm spot diameter, 400 μm rms bunch length (5 A)

Measurements and Simulations for 20-pC Bunch at 14 GeV

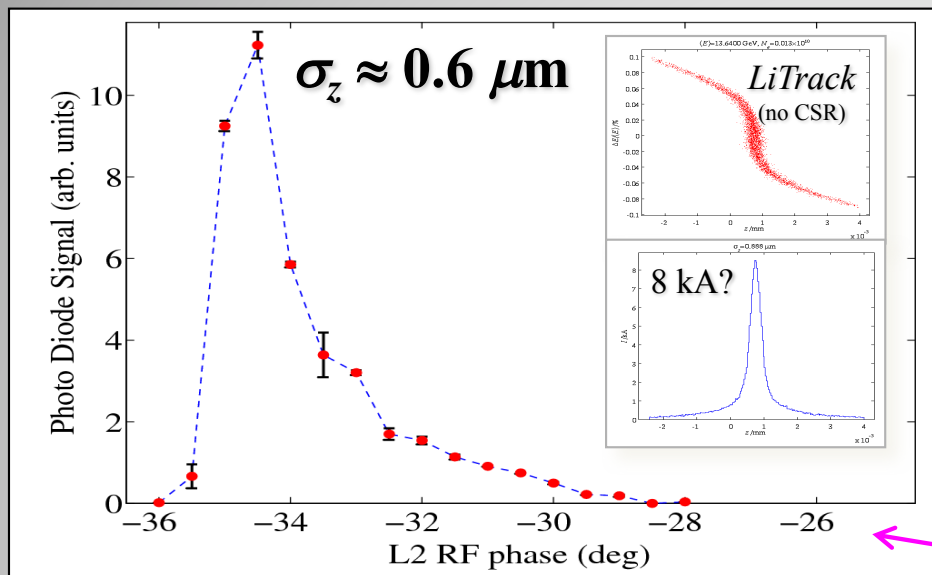
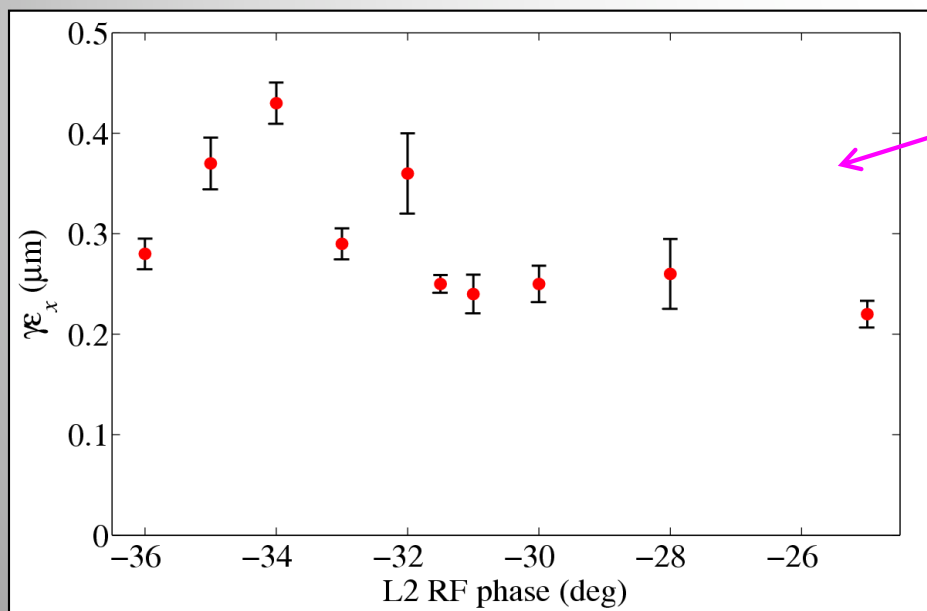
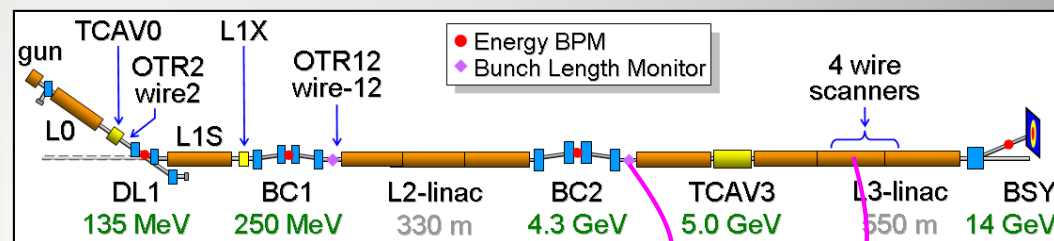
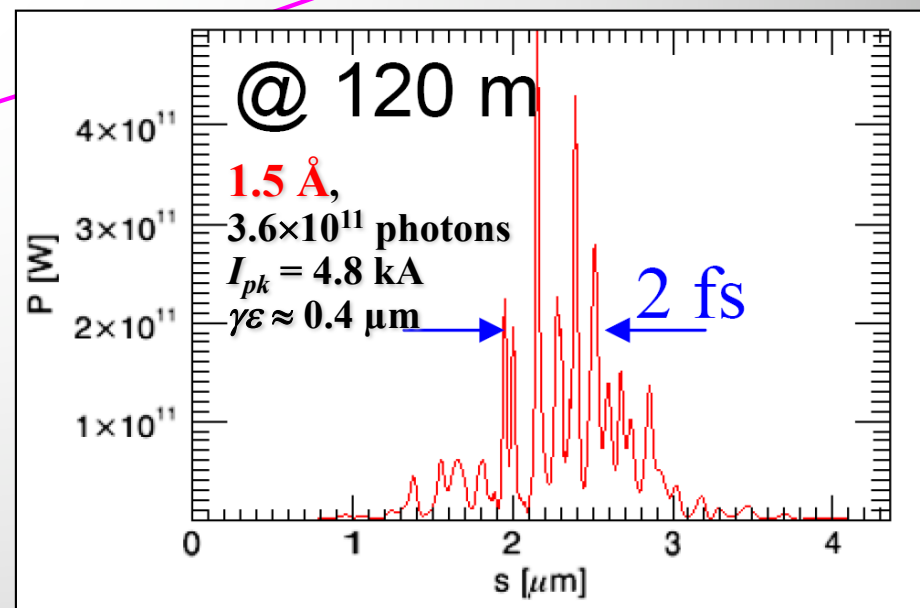


Photo-diode signal on OTR screen after BC2, best compression at L2-linac phase of -34.5 deg.



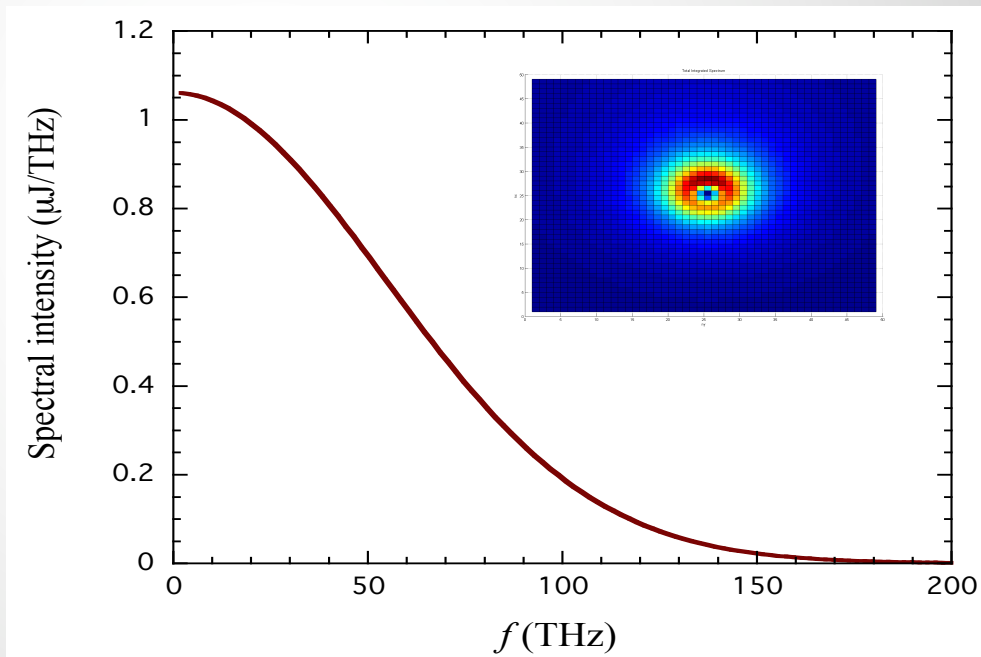
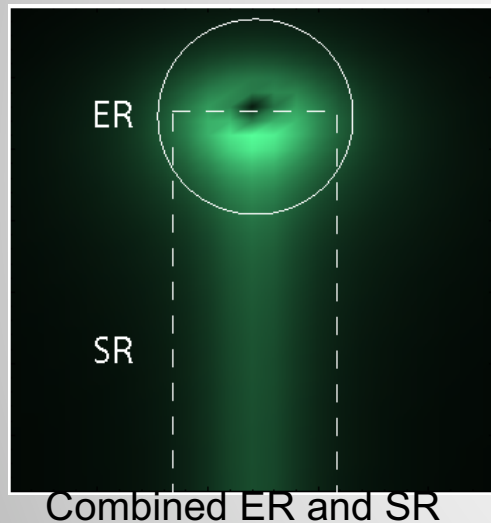
Horizontal projected emittance measured at 10 GeV



LCLS FEL simulation at 1.5 Å; not single spike.

2 fs beams at temporal measurement resolution limit

- Coherent transition radiation (destructive)
- Non-destructive: coherent edge radiation (CER)




QUINDI simulation FACET II case

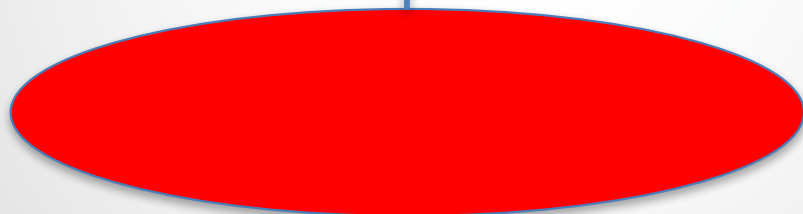
Advanced accelerator physics: focusing ultra-short beams

- 2 fs (600 nm) beam predicted to have $I_p=8$ kA
- Focus to $\sigma_r < 200$ nm (low emittance enables...)

- Surface fields

$$eE_r \approx r_e m_e c^2 I_p / ec \sigma_r$$

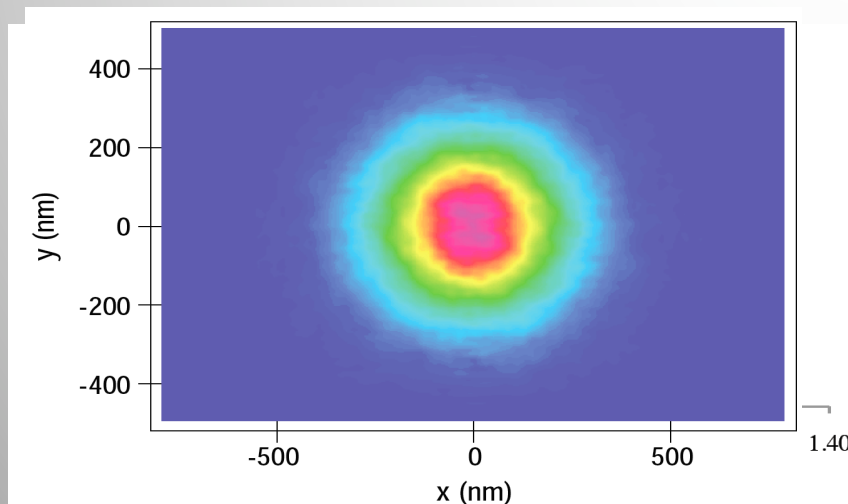

$$E_r \approx 1 \text{ TV/m!}$$



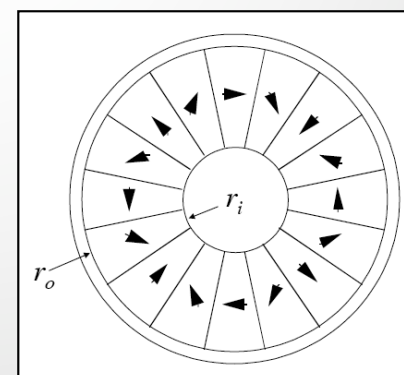
- TV/m (100 V/Å!) in fs unipolar (1/2-cycle) pulse
 - New tool for high field-matter interaction (AMO, nuclear)
 - FACET I limit ~ 100 GV/m – emittance too high!

How to focus?

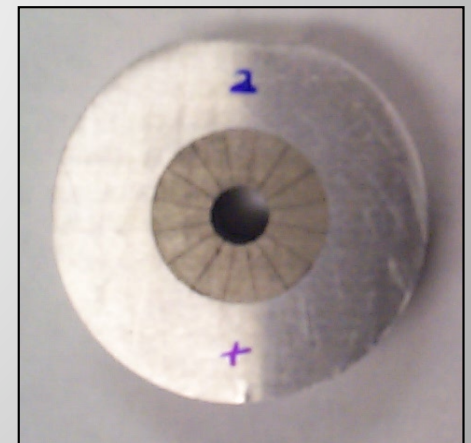
- Very short focal length final focus
- Use ultra-high field permanent magnet quads
 - mitigate chromatic aberrations
 - FF-DD-F triplet, adjust through quad placement
- Developed 570 T/m PMQ fields
 - Need slightly stronger, no problem (Pr gives $>1\text{kT/m}$)



Final beam sizes: ~ 130 nm



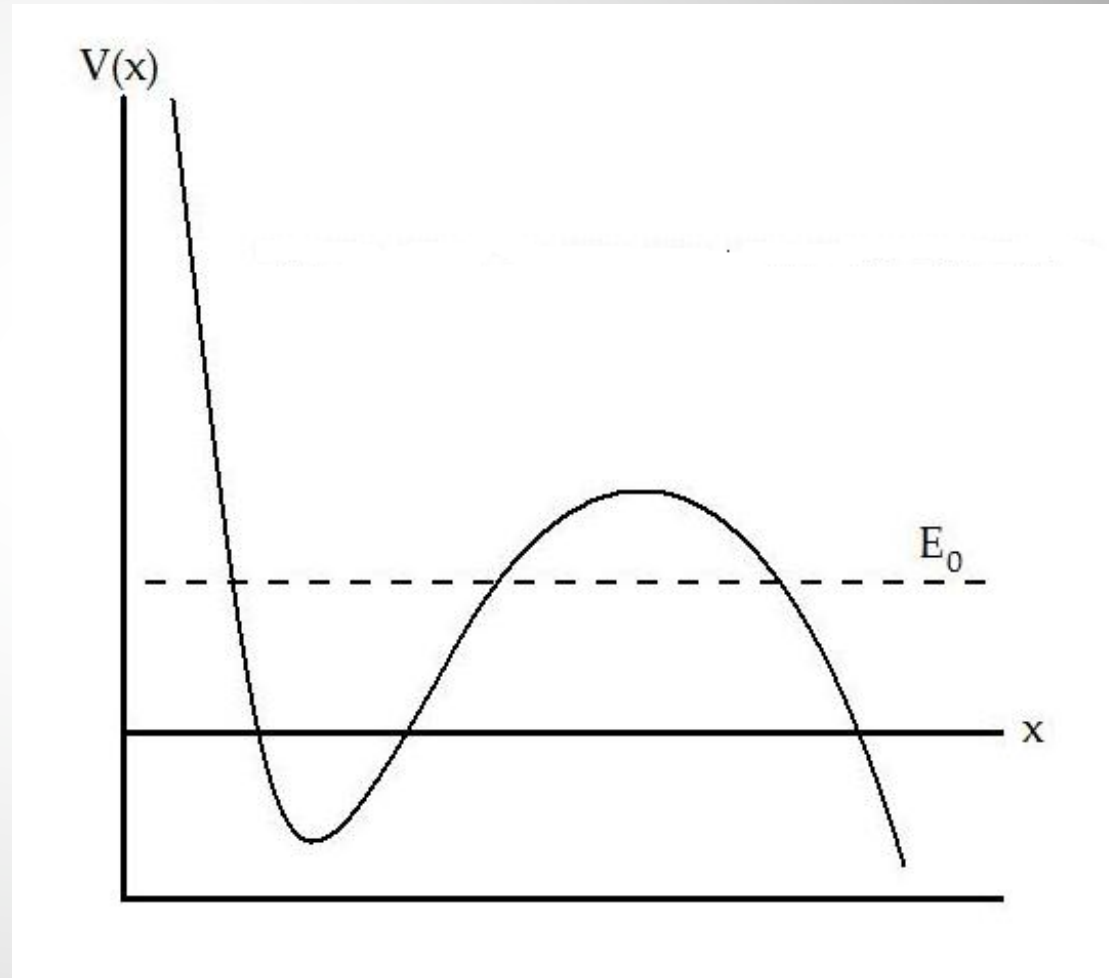
Small aperture
PMQ $B' = 570$ T/m



Halbach PMQ as built

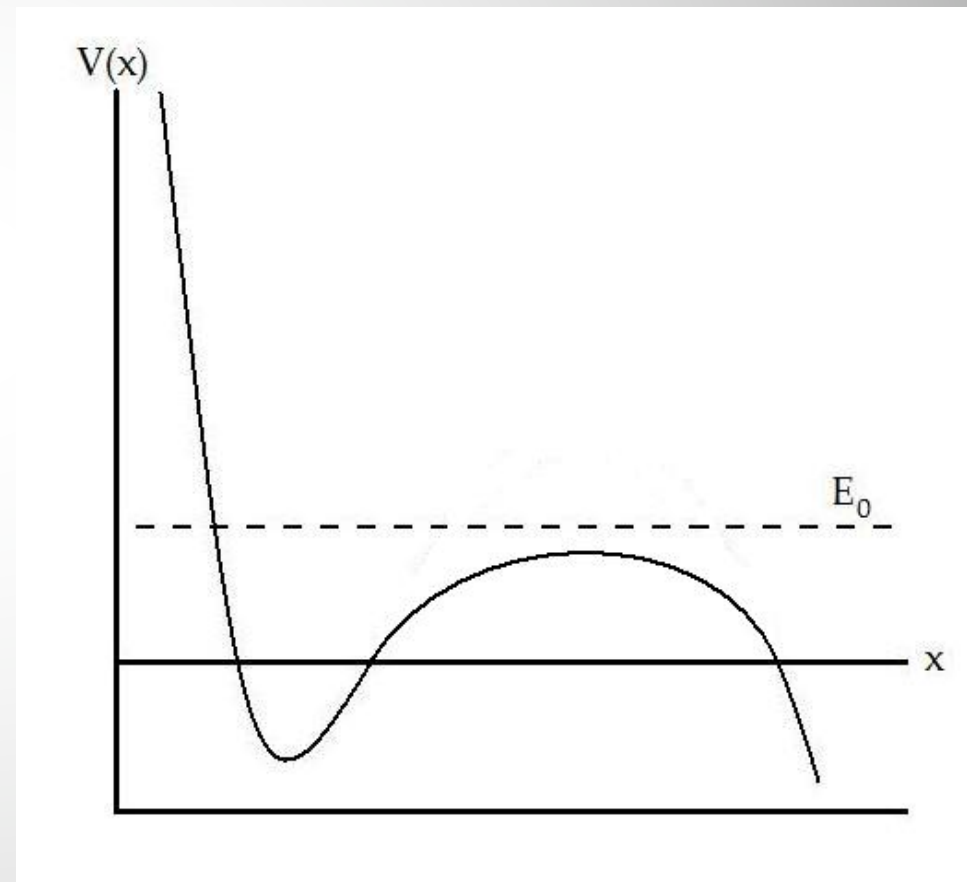
Collective Beam Field-induced Tunneling Ionization

- “Weaker” fields: tunneling
- Regime well understood
 - ADK perturbation theory
 - Developed for lasers
 - ADK-based simulation (OOPIC, Vsim)
 - Benchmarked to e-beam experiments (E167 and successors)



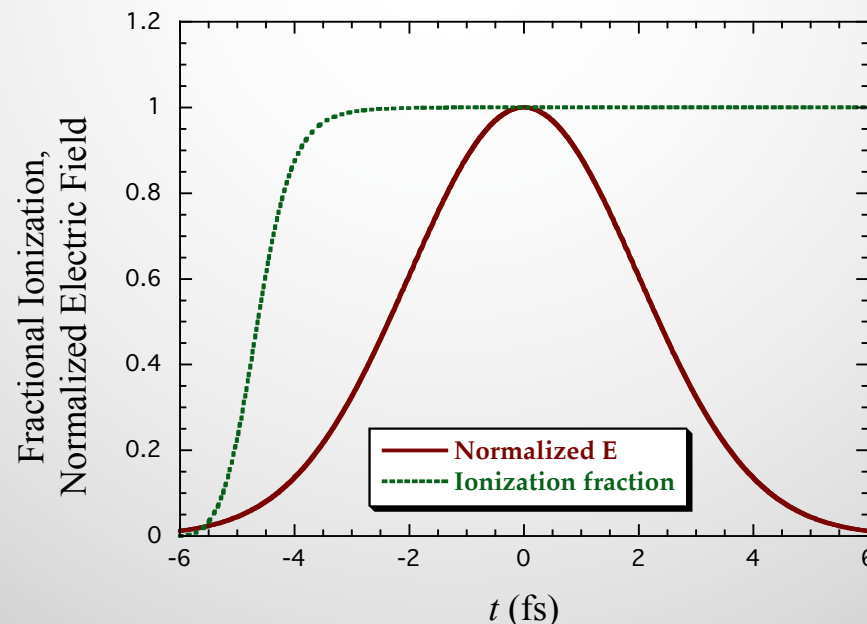
1 TV/m Reaches the Barrier Suppression Regime (BSI)

- BSI: e- classically escapes atom
 - Previously only reached experimental by lasers
 - Theory concentrates on lasers
- BSI not well understood
 - *Non-perturbative*
 - Empirical formulas
- Fundamental atomic physics tool
- Plasma wakefields...



BSI ionization occurs in 2 fs case

- Extension to unipolar field pulse
 - approach of Bauer, et al. in laser context
- BSI important above 40 GV/m, but tunneling has already been accomplished...
- For total ionization trust OOPIC

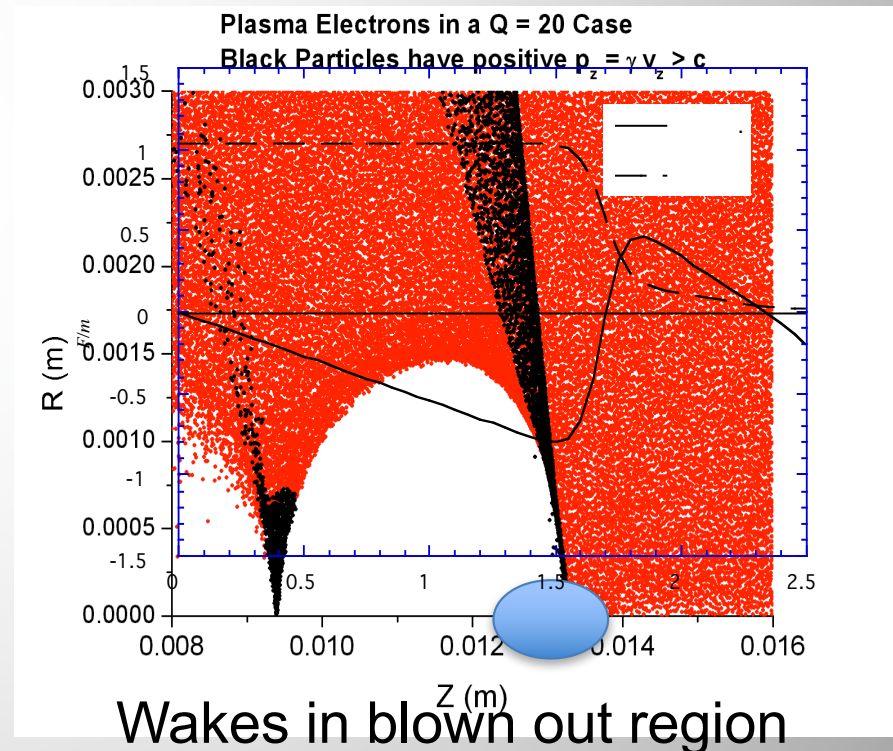


Fractional ionization due to BSI, 800 GV/m peak, 2 fs gaussian pulse

TV/m Plasma Wakefield Accelerator

- Ultra-high brightness, fs beams in plasma
- Use 20 pC LCLS beam in high n plasma
- In “blowout” regime: total rarefaction of plasma e⁻s
 - Beam denser than plasma
 - Very nonlinear plasma dynamics
 - Pure ion column focusing for e⁻s
 - EM acceleration, independent of r
 - General measure of nonlinearity:

$$\tilde{Q} \equiv \frac{N_b k_p^3}{n_0} = 4\pi k_p r_e N_b \begin{cases} \ll 1, & \text{linear regime} \\ > 1, & \text{nonlinear "blowout"} \end{cases}$$



MAGIC simulation of blowout PWFA case

Single bunch excitation at FACET II

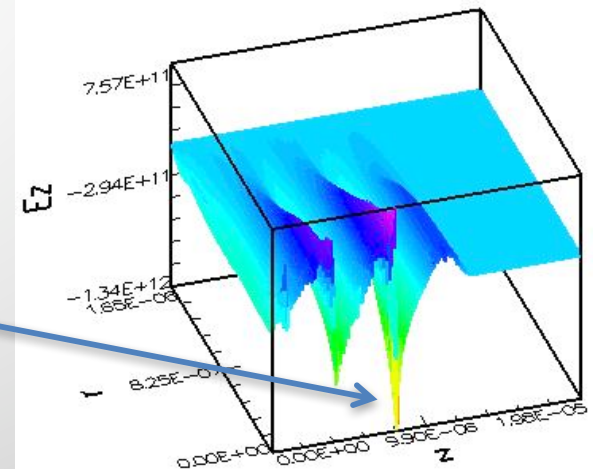
- Beam must be short and narrow compared to plasma skin depth $\sigma_r < k_p^{-1}$ $\sigma_z < k_p^{-1}$
- In this case $\tilde{Q} > 1$ implies $n_b > n_0$, blowout
- With 2 fs FACET II beam we choose $n_0 = 7 \times 10^{19} \text{ cm}^{-3}$
- For 20 pC beam, we have $\tilde{Q} = 7$

- Linear “Cerenkov” scaling

$$eE_{z,dec} \approx \frac{eE_{z,dec}^2 N_{dec} \int \frac{n(k) e^{-1} N_b}{n(k) \sigma_z^2} dk}{\sigma_z^2} \Rightarrow e^2 N_b k_p^2$$

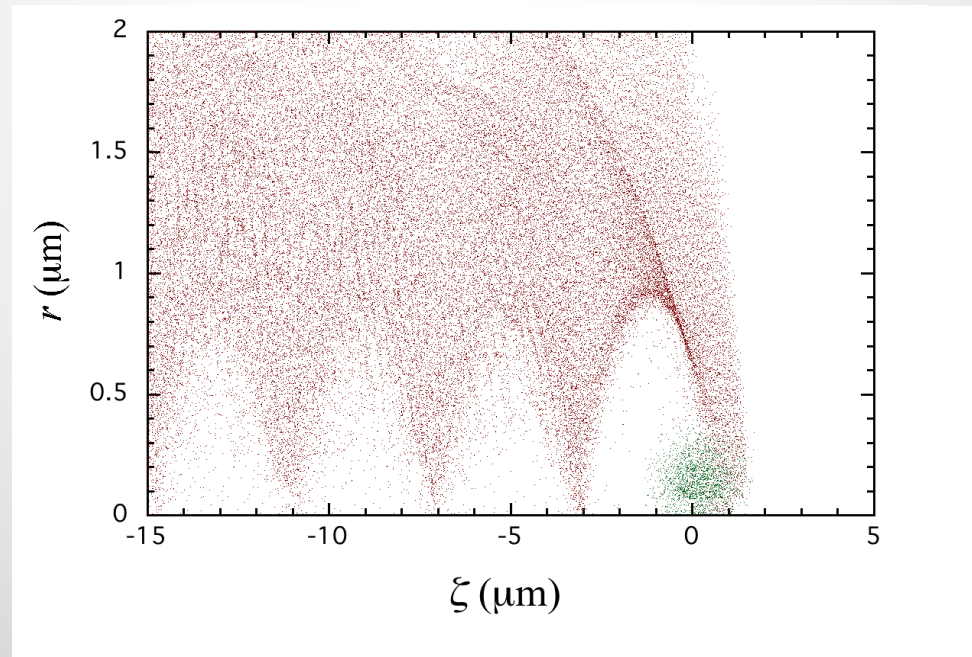
- *1 TV/m fields, converted E_r*
- Collaboration initiated (authors)

OOPIC simulation of LCLS case



Beam-field induced ionization

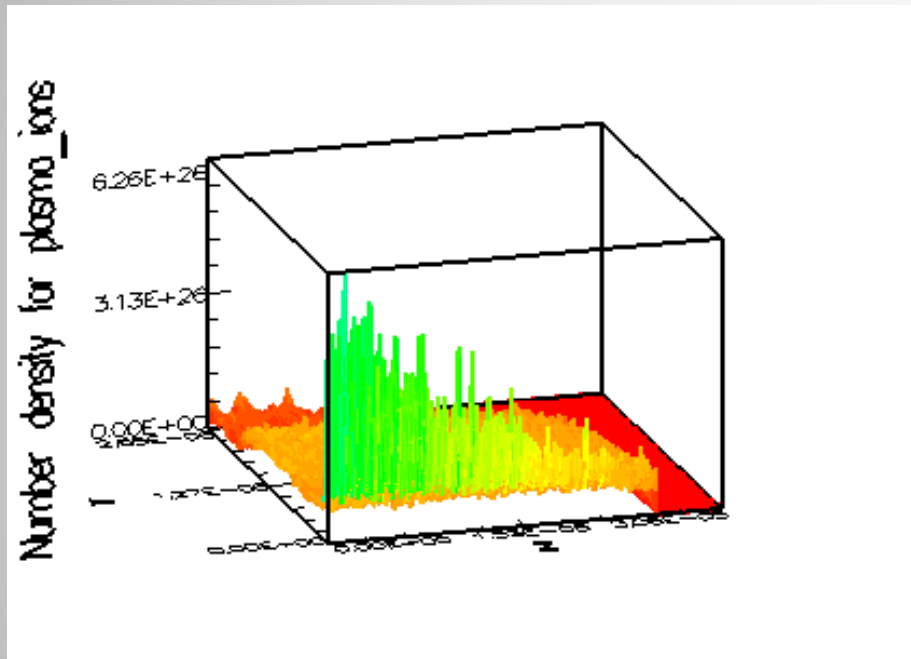
- Focus beam to < 200 nm rms
- Radial E-field $> \text{TV/m}$
- Ionization studied in Li, H gas (ADK model, which applies in beam head...)



OOP study, 315 nm, 1.5 pJ, complete ionization in Li

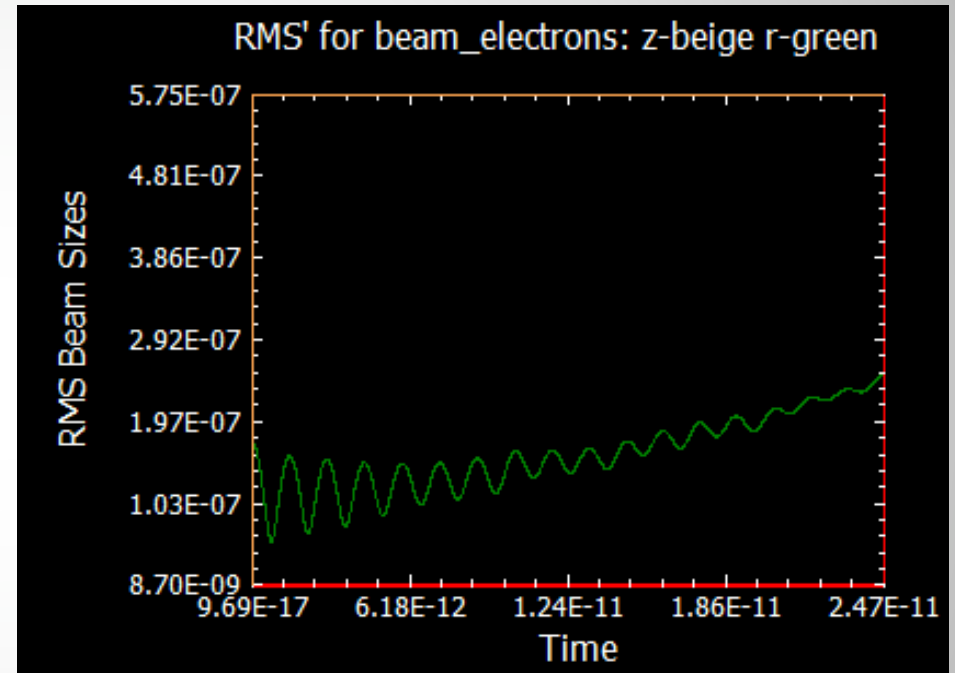
Well-focused, dense electron beam can lead to ion collapse

- Positive ions “focused” by ultra-dense e-beam fields



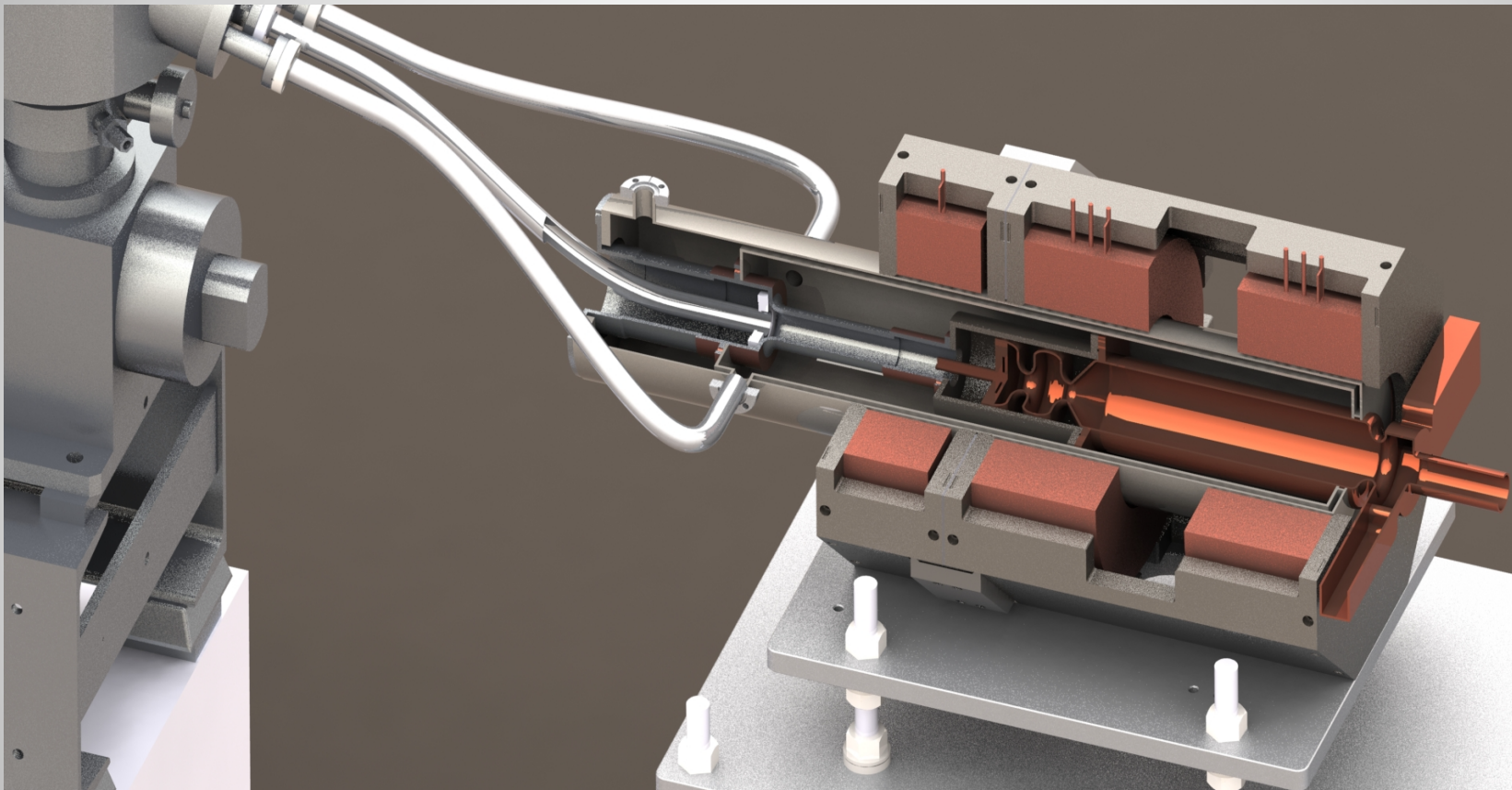
Non-uniform ion density enhancement

- Nonlinear fields, emittance growth. Bad for linear collider applications
- Detect 10-100 keV ions (hydrogen)



Beam mismatch and growth (ϵ -growth)

Increased brightness within reach



**250 MV/m peak field S-band cryogenic
gun with cryostat, focusing magnets**

Brightness at photocathode

- *Brightness* at cathode: $B_e = \frac{2I}{\varepsilon_n^2} = \frac{2J_{\max} m_e c^2}{k_B T_c}$

- In 1D limit, peak current from a pulsed photocathode is

$$J_{z,b} \approx \frac{ec\varepsilon_0}{m_e c^2} (E_0 \sin \varphi_0)^2$$

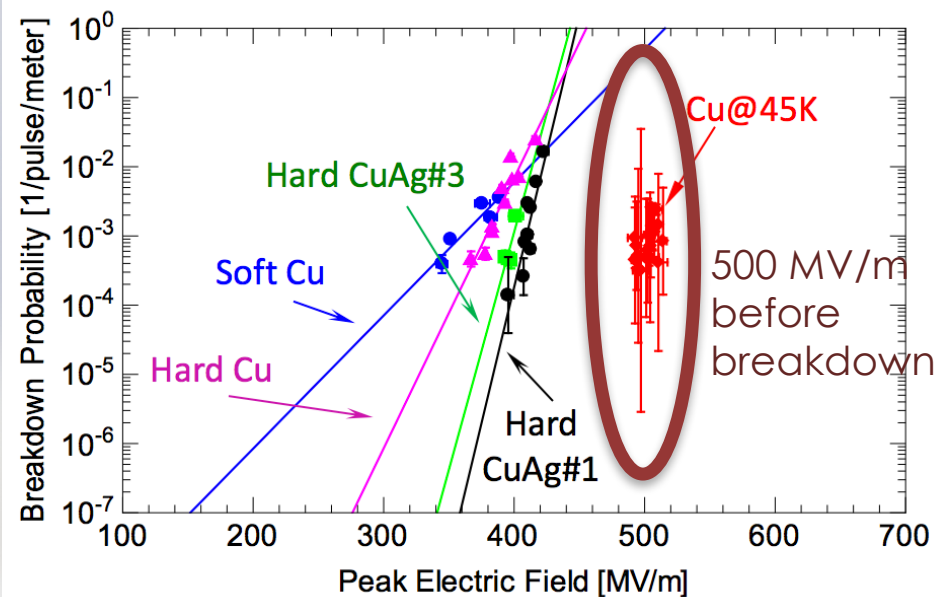
- Brightness is

$$B_{e,b} \approx \frac{2ec\varepsilon_0}{k_B T_c} (E_0 \sin \varphi_0)^2$$

- Lower emission temperature and/or...
- **Lesson: increase launch field**

Dramatically higher gradients in higher yield strength materials

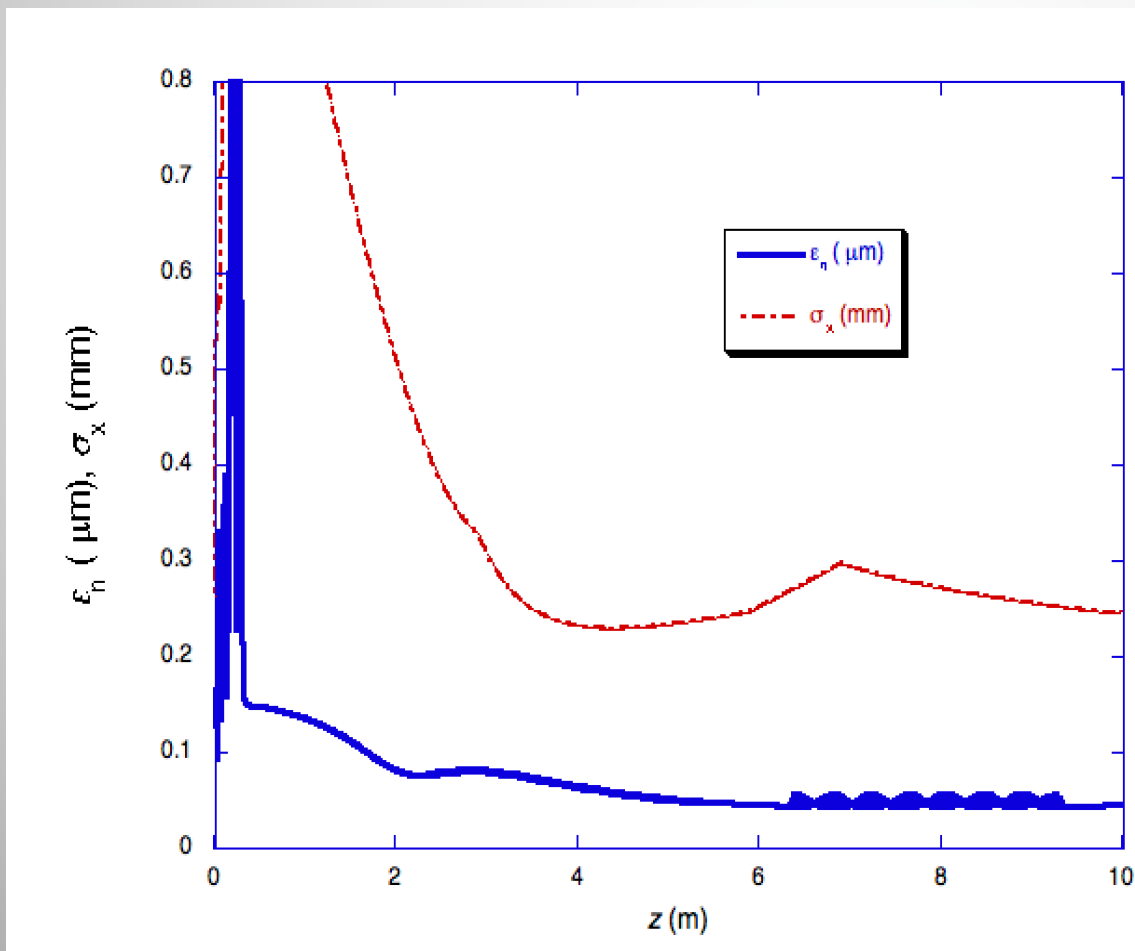
- SLAC X-band studies on hard Cu, CuAg alloy show great improvement
- ***Cryogenic structures*** give yet higher gradients, *and* lower dissipation



Game changing technique for high launch fields²²
Practical limit (dark current) ~300 MV/m presently

GPT simulation of 200 pC case

- Use long cigar-like beam (10 ps)
- Emittance $\varepsilon_n = 45$ nm-rad
 - Ten times smaller than previous example



Current $I=20$ A
Brightness a record
for this charge

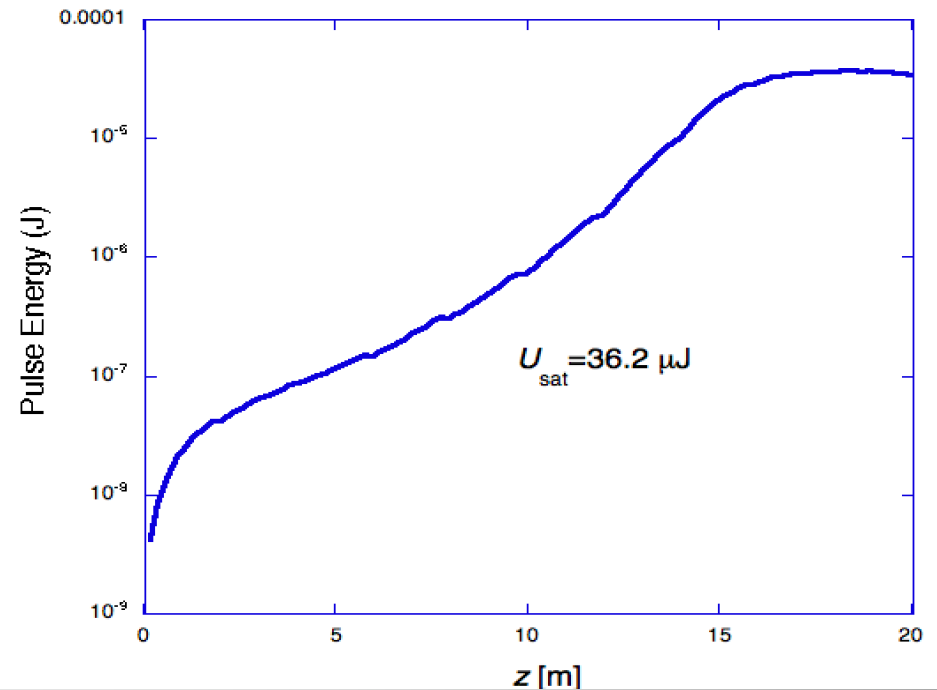
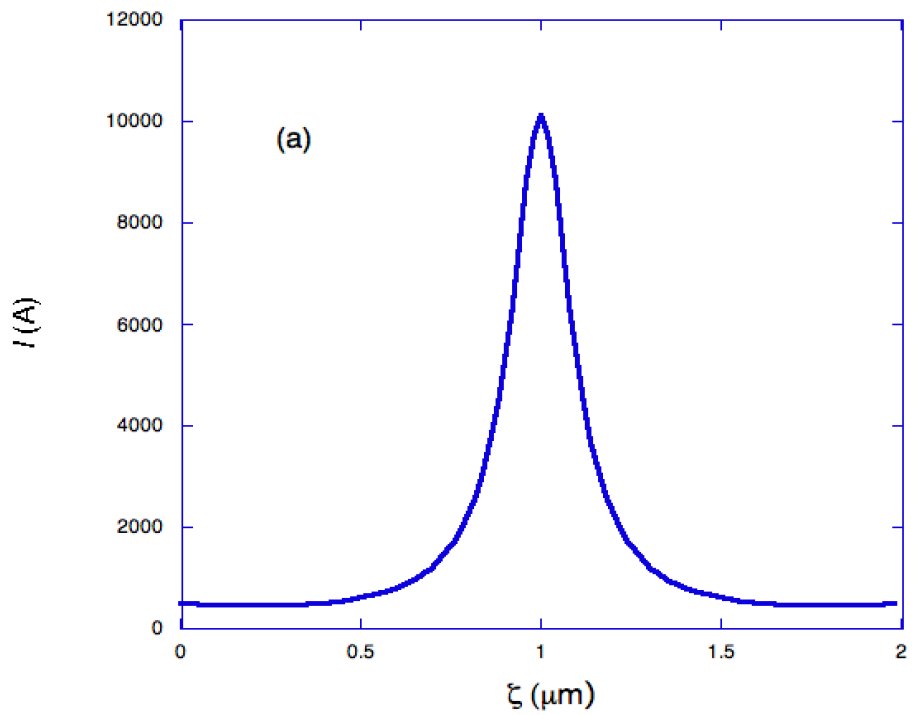
$$B_e = 2 \times 10^{16} \text{ A/m}^2$$

This is **six times**
what is available at
a *reoptimized* LCLS

**Compression is
hard. Use ESASE
microbunching**

ESASE results at 100 pC

- 100 pC (10 ps), **36 nm** emittance
- Short period *cryo-undulator*, $\lambda=9$ mm, $K=1.8$
- Operation at 14 GeV gives **80 keV X-ray**
- Saturation in <20 m, with 70 GW peak



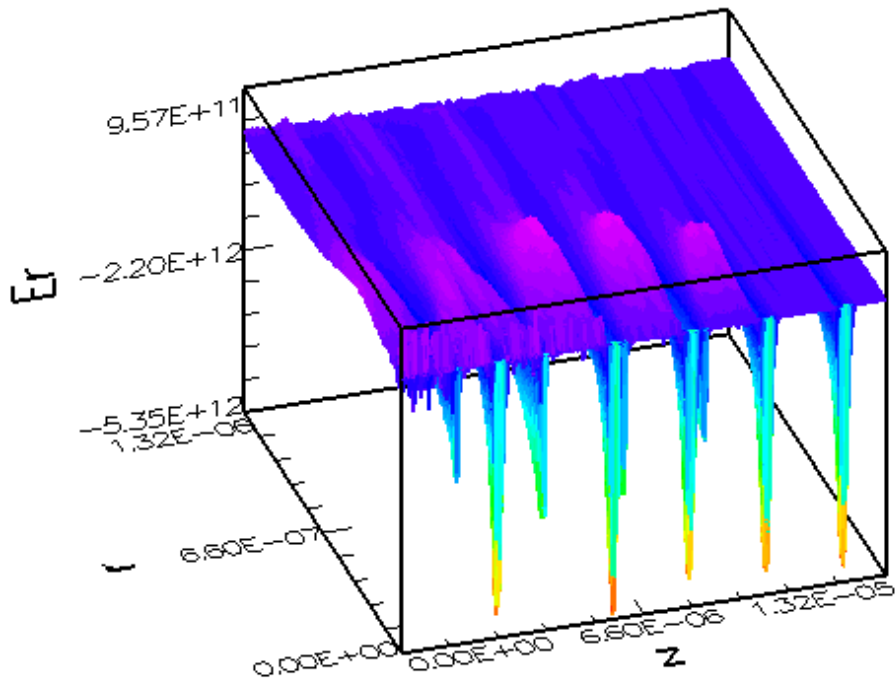
Current profile (10 kA) Energy evolution

Using 10 kA peak microbunched beam in PWFA

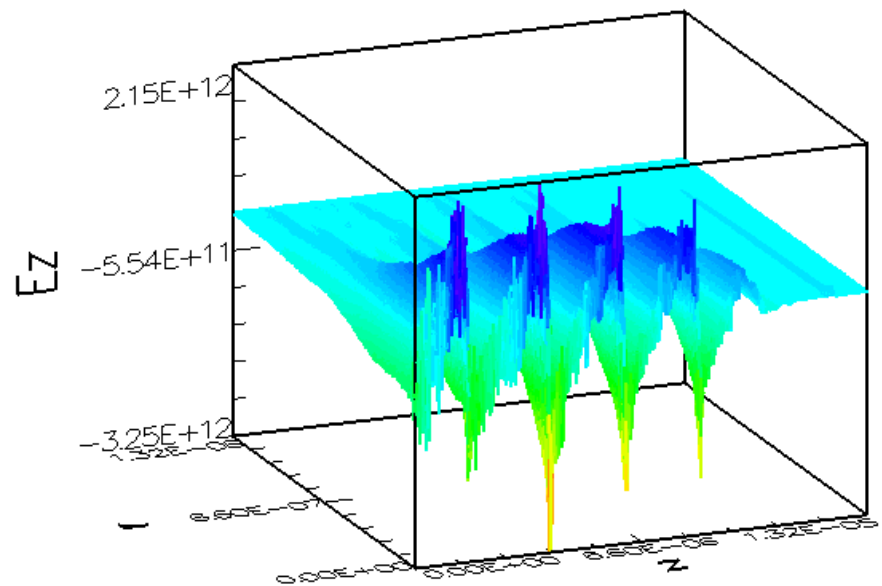
- Utilize quasi-nonlinear (QNL) regime periodic excitation (ubunching period = plasma period)

$$n_0 = 2.5 \times 10^{20} \text{ cm}^{-3}$$

- Highly focused beam (22 nm), peak $E_r > 5 \text{ TV/m}$



4 drive pulses + witness

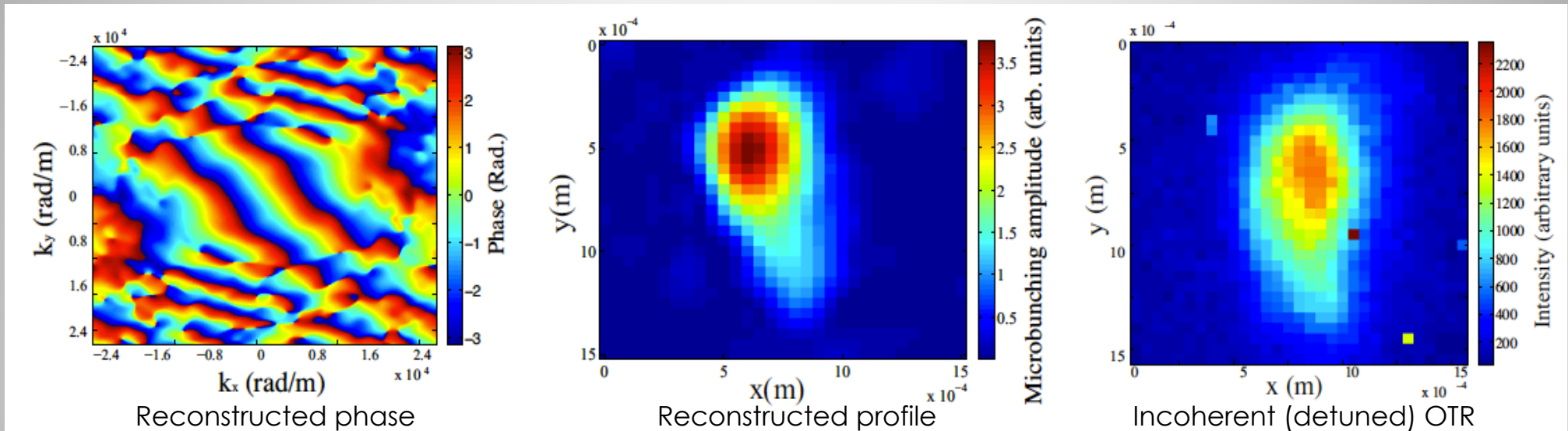


Resonant excitation to $>3 \text{ TV/m}$

Experimental implementation

- Beam focusing
 - Few-100 nm beam demands mini-beta PMQs
- Plasma section
 - $\sim 3\text{-}30$ atm gas jet, with BSI. Start with tenuous gas
 - Length ~ 0.5 mm gives $\sim \text{GeV } \Delta E$, “perturbative”
- Beam diagnostics in entirely new regime
 - Longitudinal: coherent edge/transition radiation
 - Transverse:
 - Ionization, appearance intensity
 - coherence

Sub- μm beam transverse diagnosis



Coherent transition radiation imaging reconstruction
expt., A. Marinelli et al., *PRL* 110, 094802 (2013)

- Measure sub- μm beam sizes with *coherent imaging* (borrowed from XFEL). Coherent information down to 100 nm?

Conclusions

- Attosecond e-beams can be reached at low Q
- Greatly enhanced beam brightness
 - Single spike, compact FELs
 - New high field sources, enhance wavelength range
- Frontier regime for beams; coherent optical radiation, ionization, new diagnostics
- Enables new frontiers:
 - Extreme plasma wakefield accelerators
 - TeV/m at 1 atm
 - Resonantly driven optical/IR wavelengths (beam is easier, better!)
 - Ultra-high field atomic-physics – TV/m unipolar field