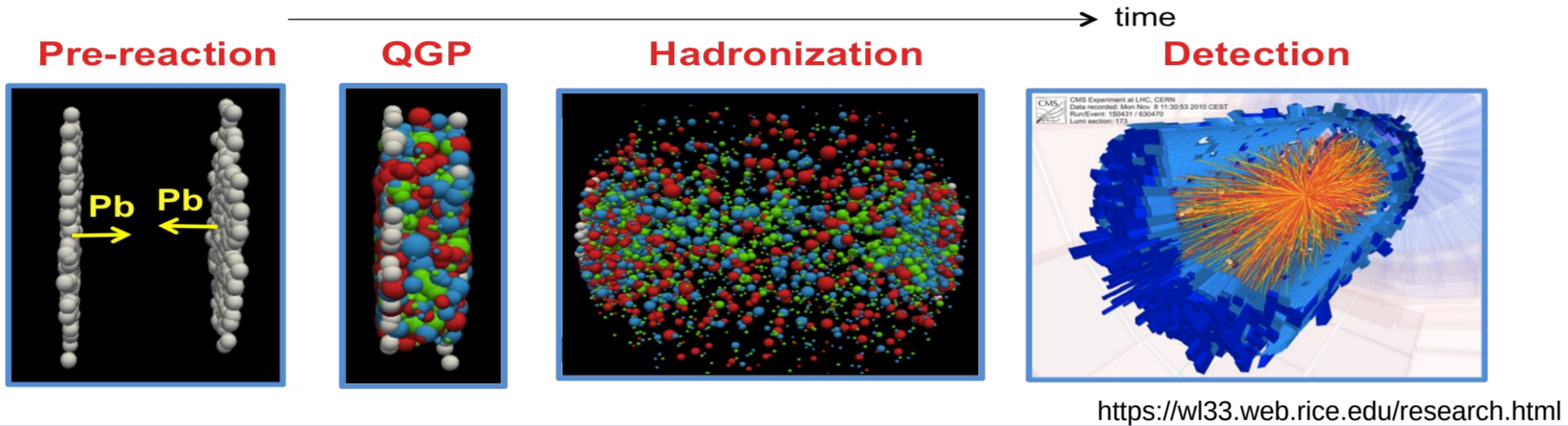
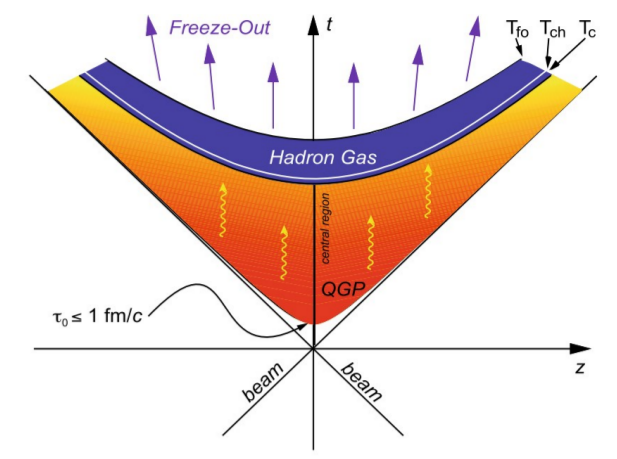


Motivation

Studying the QCD matter produced under extreme condition of temperature and density called Quark Gluon Plasma (QGP) is among the important goal of heavy-ion collision experiments. QGP state is being created for a very short interval of time (~10⁻²² s) so we cannot directly probe this state. Hence we utilize kinematic data of final state particles produced in heavy-ion collision in order to study the dynamics of QGP. Transverse momentum (p_T) spectra is one such kinematic variable that gives us information about the thermodynamical as well as hydrodynamical properties of the system produced in heavy-ion collision. We have developed a unified formalism to study full range of p_T-spectra including both soft as well as hard part using a single distribution function.

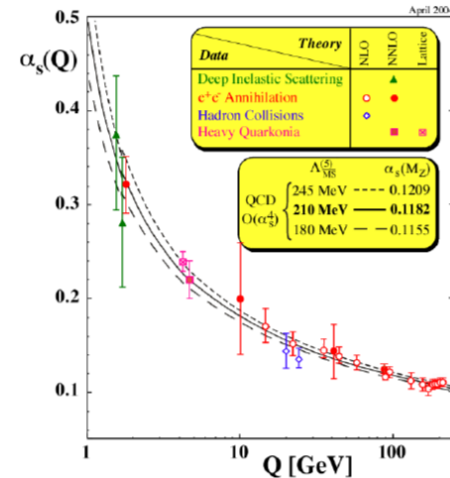


Quantum Chromo Dynamics (QCD)

QCD is the field theoretical framework which governs the strong interaction between quarks and gluons.

Theoretical models for p_T spectra

- Due to the asymptotic freedom and very nature of QCD coupling constant, it is difficult to apply perturbative QCD at low energy because of high coupling strength.
- To overcome this issue, we resort to phenomenological models with most common being the statistical approach to explain low-p_T part of the spectra whereas we have a well defined perturbative QCD based power-law form of distribution function for high-p_T region corresponding to particles produced in hard processes.



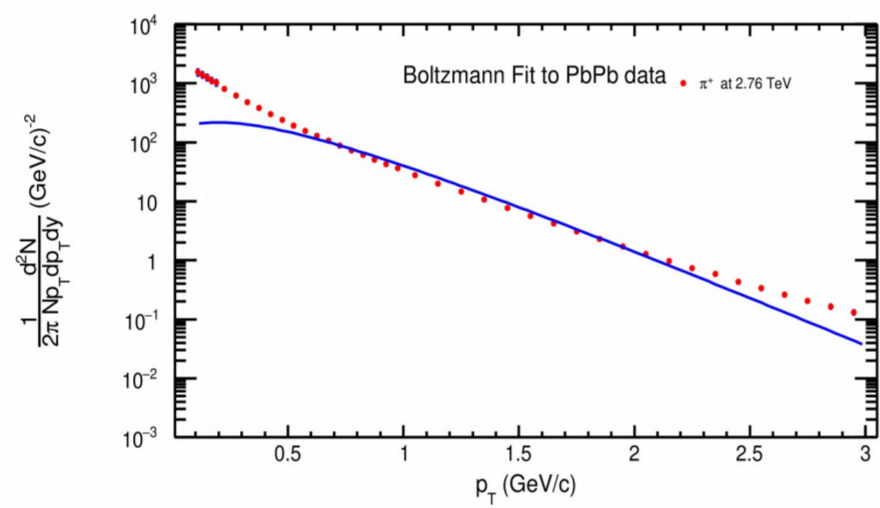
Conventional Approach to study p_T spectra

- Considering the particles produced in heavy-ion collision to be of thermal origin. Most natural choice to explain energy spectra is Boltzmann distribution.
- For Boltzmann distribution, p_T spectra is given as

$$\frac{1}{2\pi p_T} \frac{d^2 N}{dp_T dy} = \frac{g V m_T}{(2\pi)^3} \exp\left(-\frac{m_T}{T}\right)$$

- Here m_T is the transverse mass given as

$$m_T = \sqrt{m^2 + p_T^2}$$

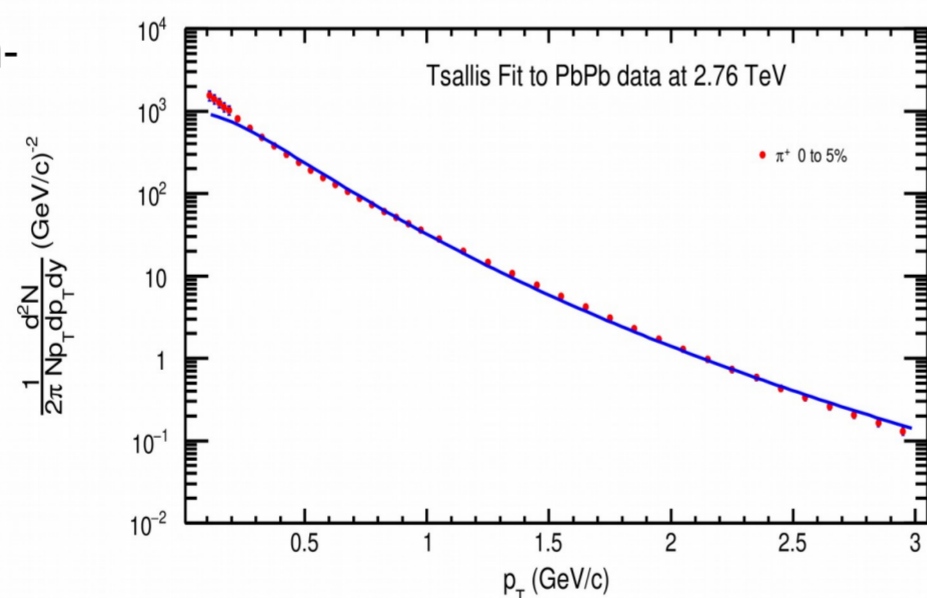


- In the graph above, we have fitted p_T-spectra of positive pions produced in 2.76 TeV Pb-Pb collision with the Boltzmann distribution.
- We observe that Boltzmann distribution deviates significantly from data beyond certain p_T range.
- In order to overcome this problem, Tsallis statistics has been introduced in high energy physics.

- Tsallis statistics [2] is a generalised Boltzmann-Gibbs statistics which also takes into account non-extensivity in the system.
- Non-extensivity can arise in strongly coupled system.

$$\frac{1}{2\pi p_T} \frac{d^2 N}{dp_T dy} = \frac{g V m_T}{(2\pi)^3} \left[1 + (q-1) \frac{m_T}{T}\right]^{-\frac{q}{q-1}}$$

- Non-extensivity parameter "q" takes care of deviation from thermal equilibrium.
- Tsallis distribution deviates from data at high p_T region corresponding to hard scattering.



Pearson Distribution

- Hard scattering part of p_T spectra is governed by power law form:

$$f(p_T) = \frac{1}{N} \frac{dN}{dp_T} = A p_T \left(1 + \frac{p_T}{p_0}\right)^{-n}$$

- Pearson distribution [3] is a generalised form of many probability distribution functions like gaussian, exponential, gamma distributions etc.
- It is given in form of differential equation:

$$\frac{1}{p(x)} \frac{dp(x)}{dx} + \frac{a+x}{b_0 + b_1 x + b_2 x^2} = 0$$

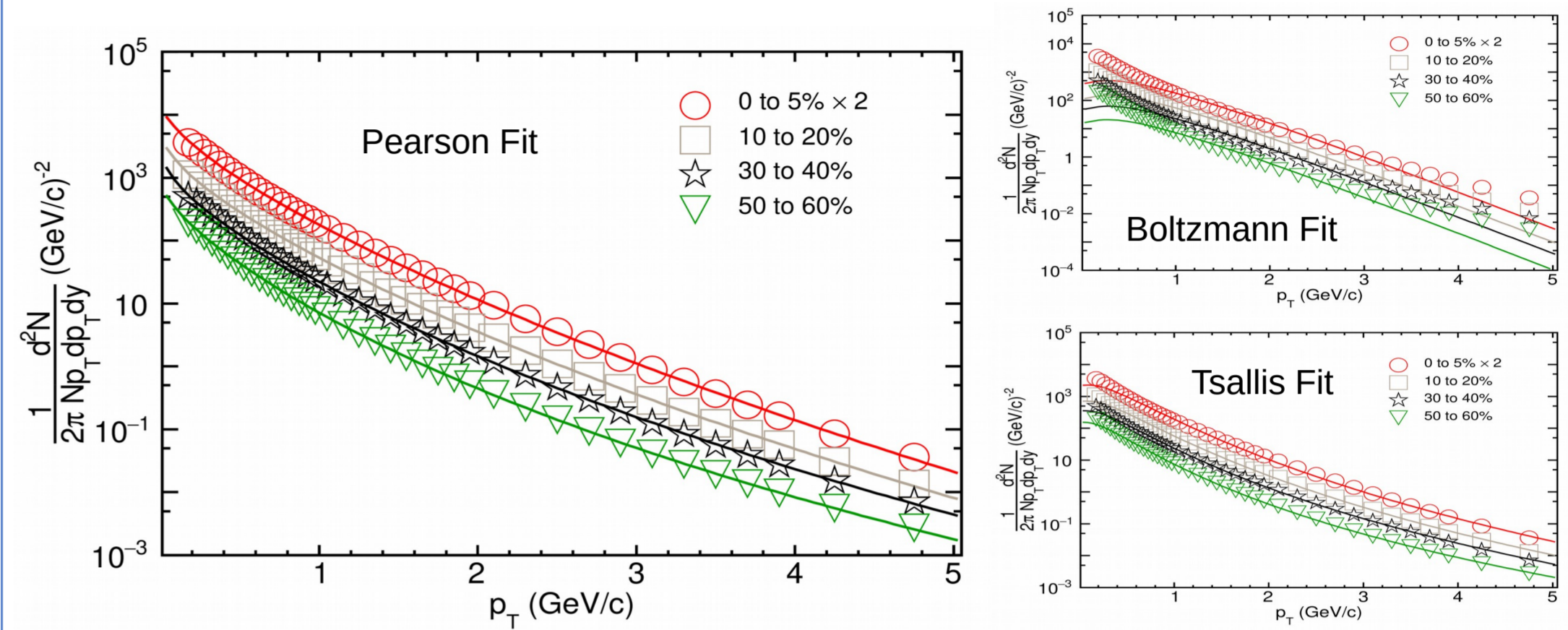
- Parameters a, b₀, b₁, b₂ are related to first four moments of a distribution.
- We have modified solution of this distribution function by substituting physics parameters to give transverse momentum spectra [1]

$$\frac{1}{2\pi p_T} \frac{d^2 N}{dp_T dy} = B \left(1 + \frac{p_T}{p_0}\right)^{-n} \left(1 + (q-1) \frac{p_T}{T}\right)^{-\frac{q}{q-1}}$$

Hard-processes Soft-processes

Here parameter values are such that the "soft" part decay very quickly and hence giving way for dominance of "hard" part after certain p_T value.

Results



Plots of p_T-spectra of charged hadrons produced in 2.76 TeV Pb-Pb collision [4] at four different centralities fitted with different distribution function.

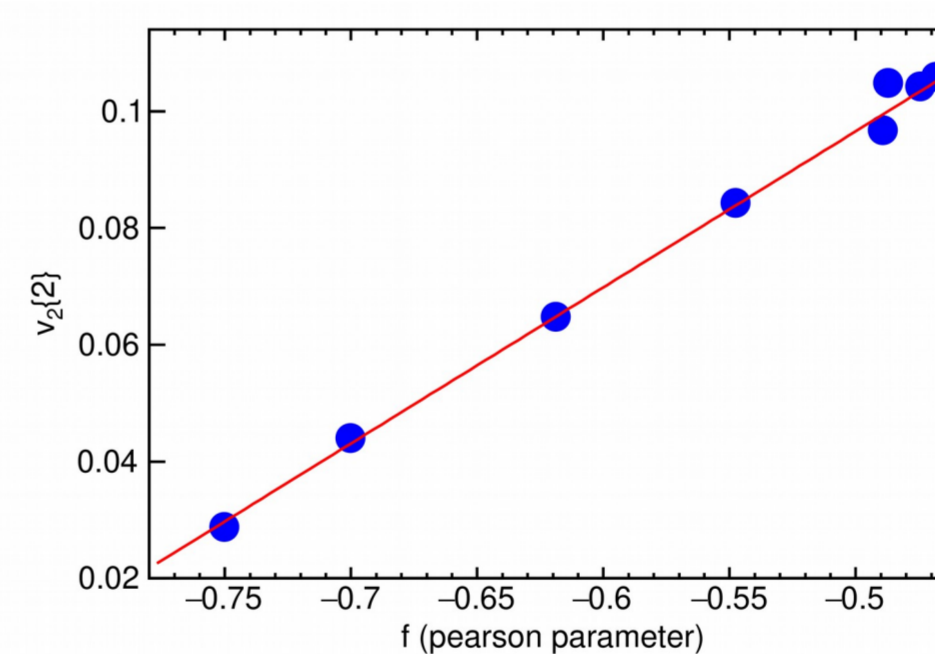
Flow Analysis

- Flow corresponds to the azimuthal anisotropy in distribution of particle produced in heavy ion collision.

$$E \frac{d^3 N}{dp^3} = \frac{1}{p_T} \frac{d^2 N}{dp_T dy} \frac{N}{2\pi} \left[1 + 2 \sum_n v_n \cos\{n(\phi - \psi)\}\right]$$

- Here, v_n is the nth order flow coefficient.

Final state momentum anisotropy v_n is correlated to initial spatial eccentricity ε_n



We observed a linear relationship between one of the Pearson fit parameter with elliptic flow coefficient v₂{2} obtained from Ref [5].

Summary

- Tsallis distribution deviates from data as we move towards higher p_T region.
- We developed a generalized approach to study both low as well as high-p_T regions of the spectra.
- We also observe that there is a linear relationship between one of the fitted parameters and elliptic flow coefficient.

Acknowledgement

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